

# FRLDM FISSION YIELDS FOR $r$ -PROCESS NUCLEOSYNTHESIS



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*Workshop on nuclear astrophysics*

Beihang University

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FIRE Collaboration

Fission In R-process Elements

# TO UNDERSTAND THE FORMATION OF THE ELEMENTS

Requires deep knowledge of a range of fields, including:

The theoretical **modeling of astrophysical environments**

**Multi-messenger observations** (gravitational waves, EM waves, etc.)

**Nuclear theory predictions** for exotic nuclei

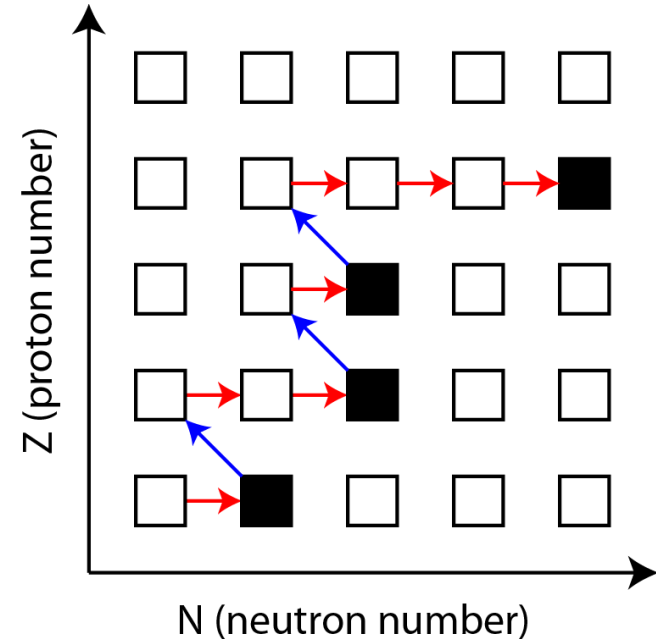
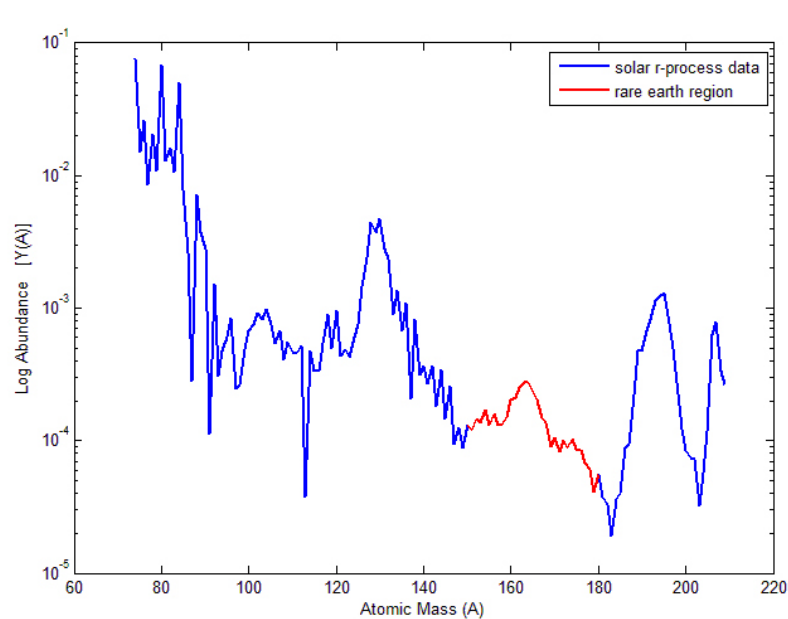
**Precision experiments** to constrain nuclear theory

Data and observations are **limited**

We must be clever when deciphering what is going on with nucleosynthesis...



# WHAT IS THE $r$ -PROCESS?



Rapid neutron capture that occurs in astrophysical environments allowing for the production of **heavy elements**

**Neutron captures** are initially much faster than  $\beta$ -decays

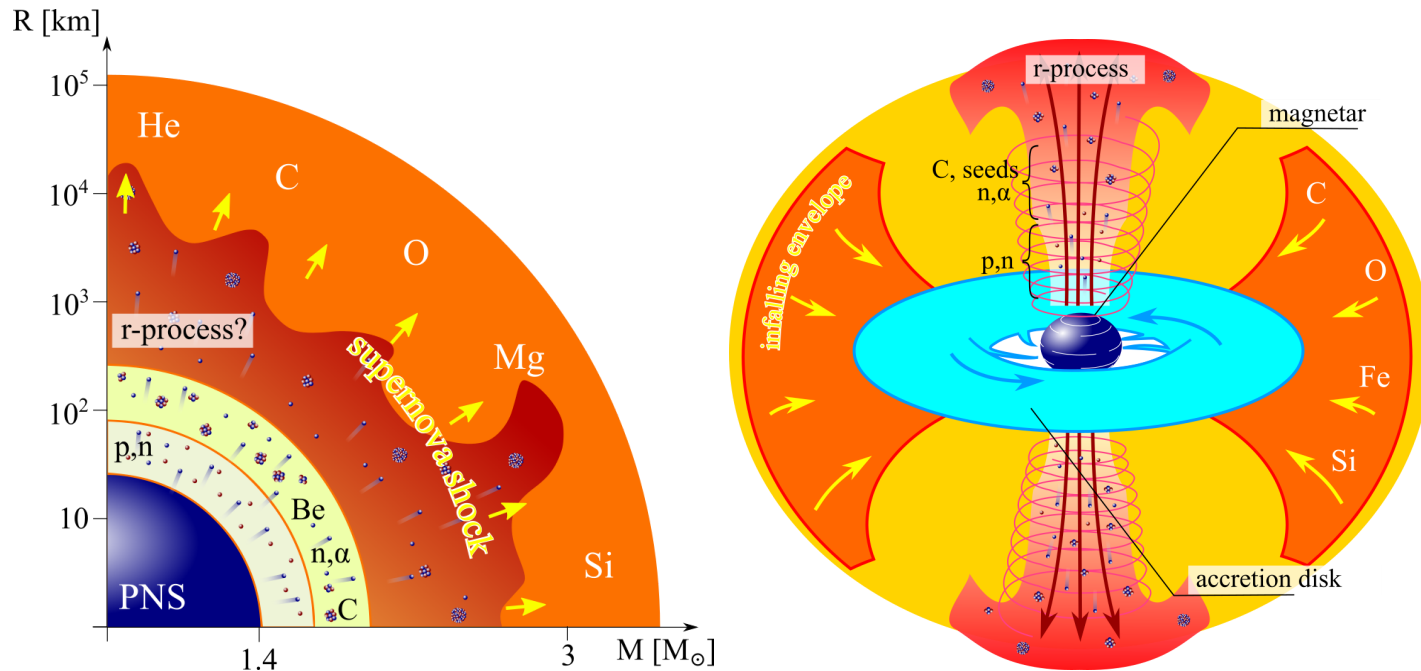
Relative slowdown in the nuclear flow (right) produces peak structures in the observed abundances (left)

Astrophysical environment must produce a lot of free neutrons in order for this process to proceed



# WHERE CAN THE $r$ -PROCESS OCCUR?

One possibility is in (rare?) supernovae



For standard supernovae (left) neutrino physics still needs to be well understood

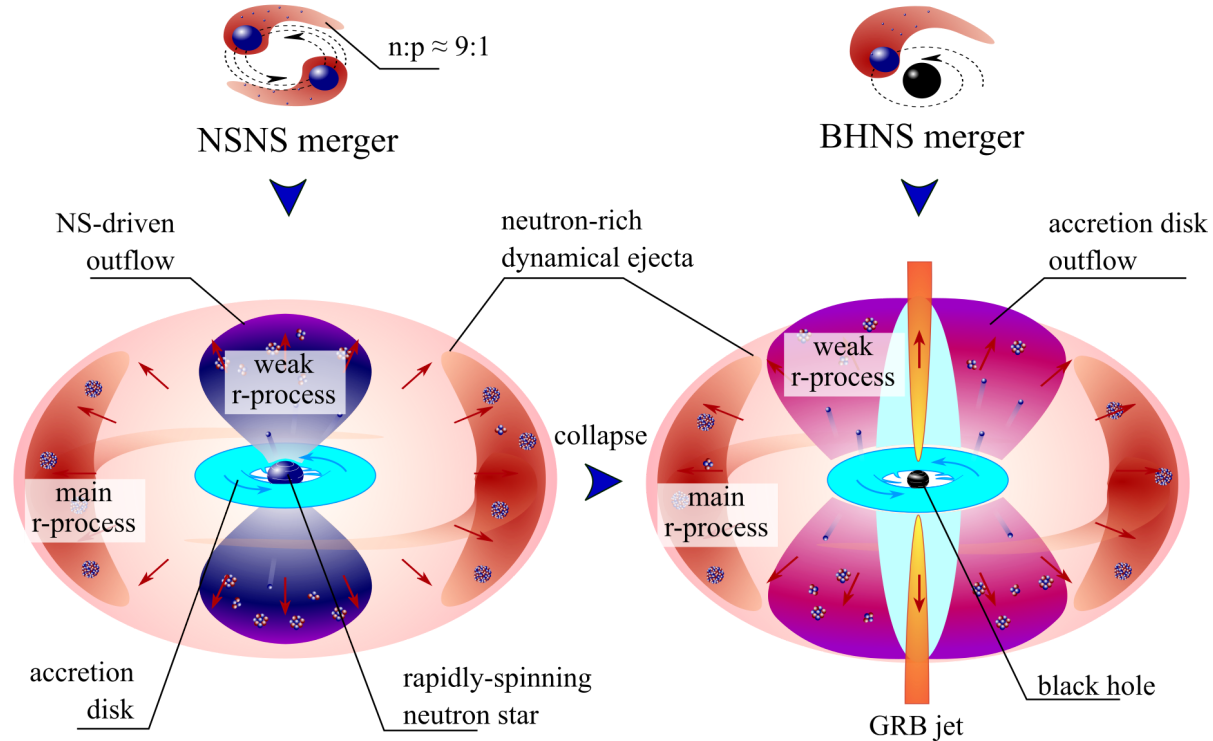
Jets in magnetorotational driven supernovae (right) may also provide the necessary conditions

Another option is the disk winds of collapsars - black hole forms after core collapse of a rapidly rotating star



# WHERE CAN THE *r*-PROCESS OCCUR?

Another possibility is in compact object mergers

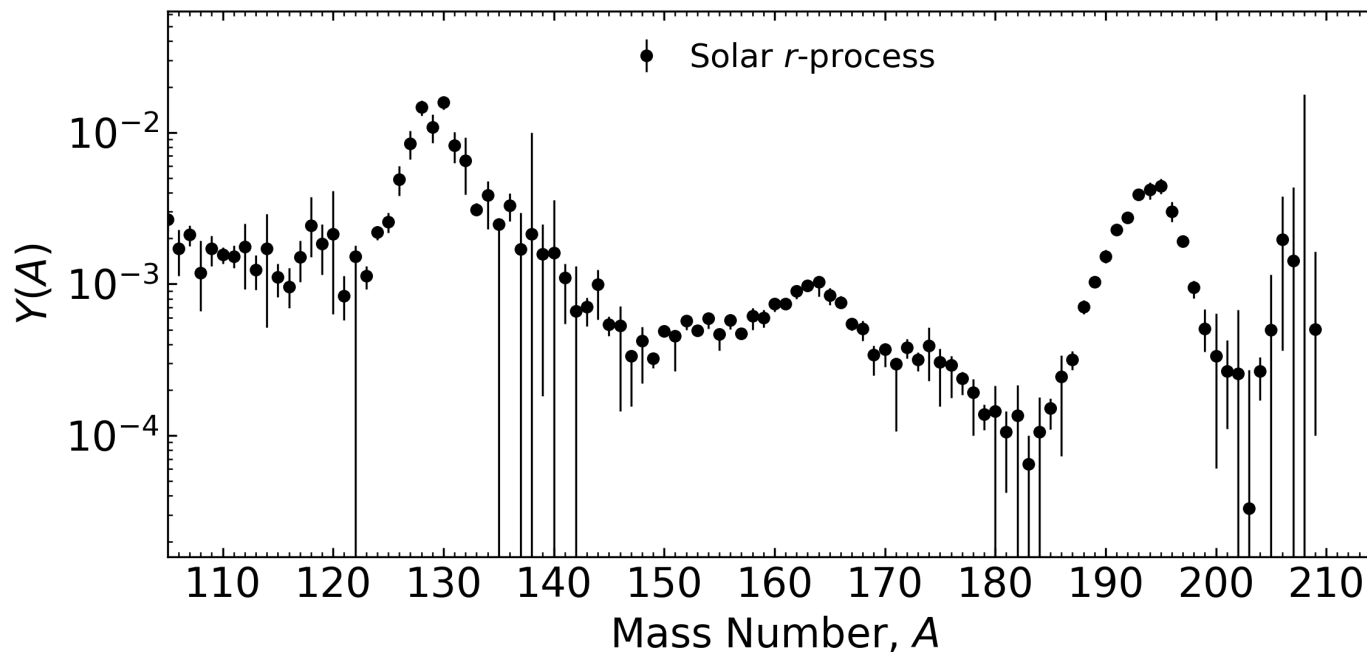


A binary merger of neutron stars is an exciting possibility (some indirect evidence exists)

Another option is in the disk of a black hole neutron star binary

# WHEN WE MODEL NUCLEOSYNTHESIS

We want to describe the abundances observed in nature



But there is uncertainty in:

The astrophysical conditions (large variations in current simulations)

The nuclear physics inputs (1000's of unknown species / properties)

# INPUTS FROM NUCLEAR PHYSICS

**1st order:** masses,  $\beta$ -decay rates, reaction rates & branching ratios





# $r$ -PROCESS CALCULATION

**nuclear physics inputs**

( $S_n$ ,  $\beta$ -rates,  $n$ -cap rates, ... )

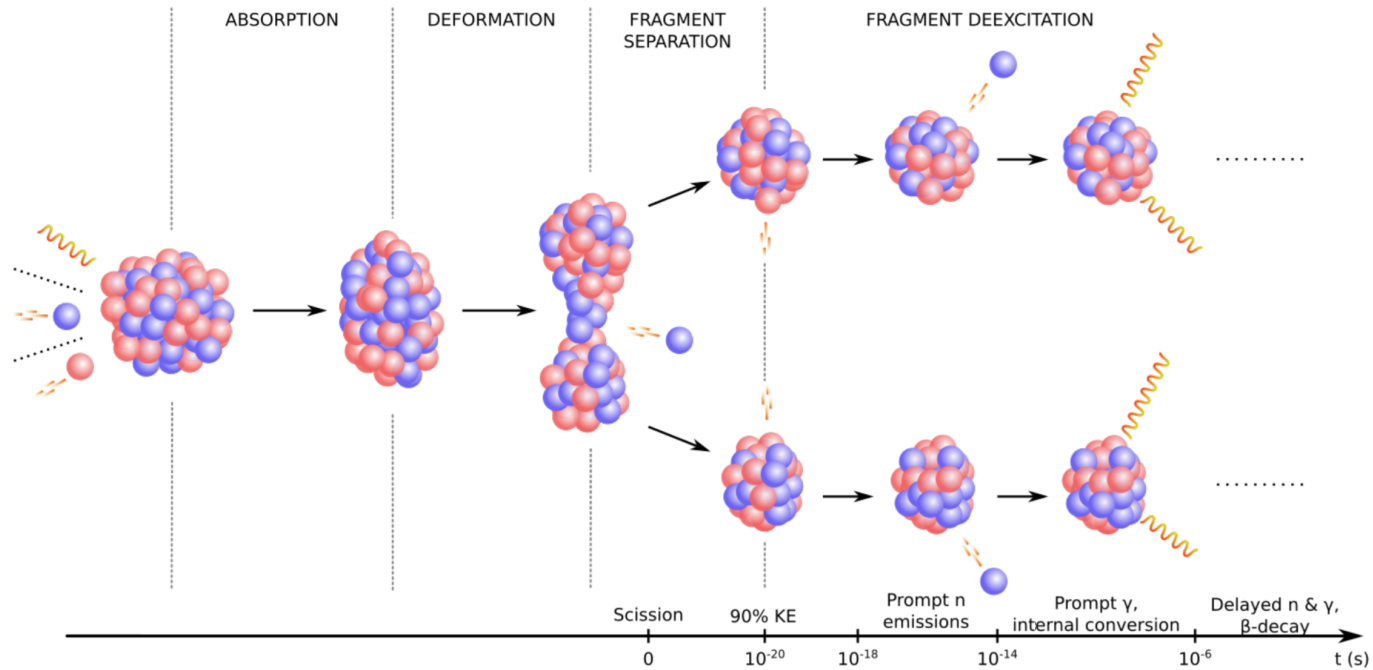


**thermodynamic conditions**

(temperature, density, ... )

PRISM: Portable Routines for Integrated nucleoSynthesis Modeling

# NUCLEAR FISSION IN A NUTSHELL



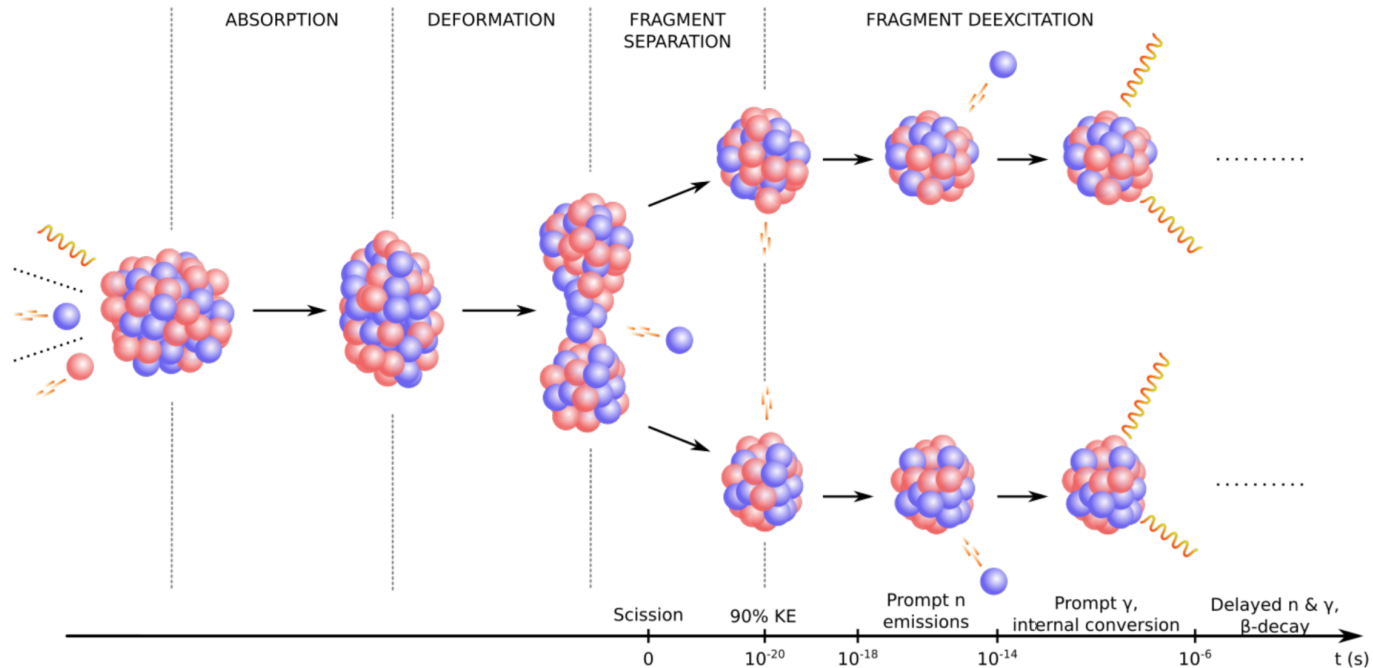
## The fission process:

A heavy nucleus splits into two lighter fragments

Subsequent particle emission and decays then occur

Many events gives rise to fission yield

# NUCLEAR FISSION FOR THE $r$ -PROCESS



## Influence on the $r$ -process:

Fission **rates** and **branching** determine re-cycling (robustness)

Fragment **yields** place material at lower mass number; barriers determine hot spots

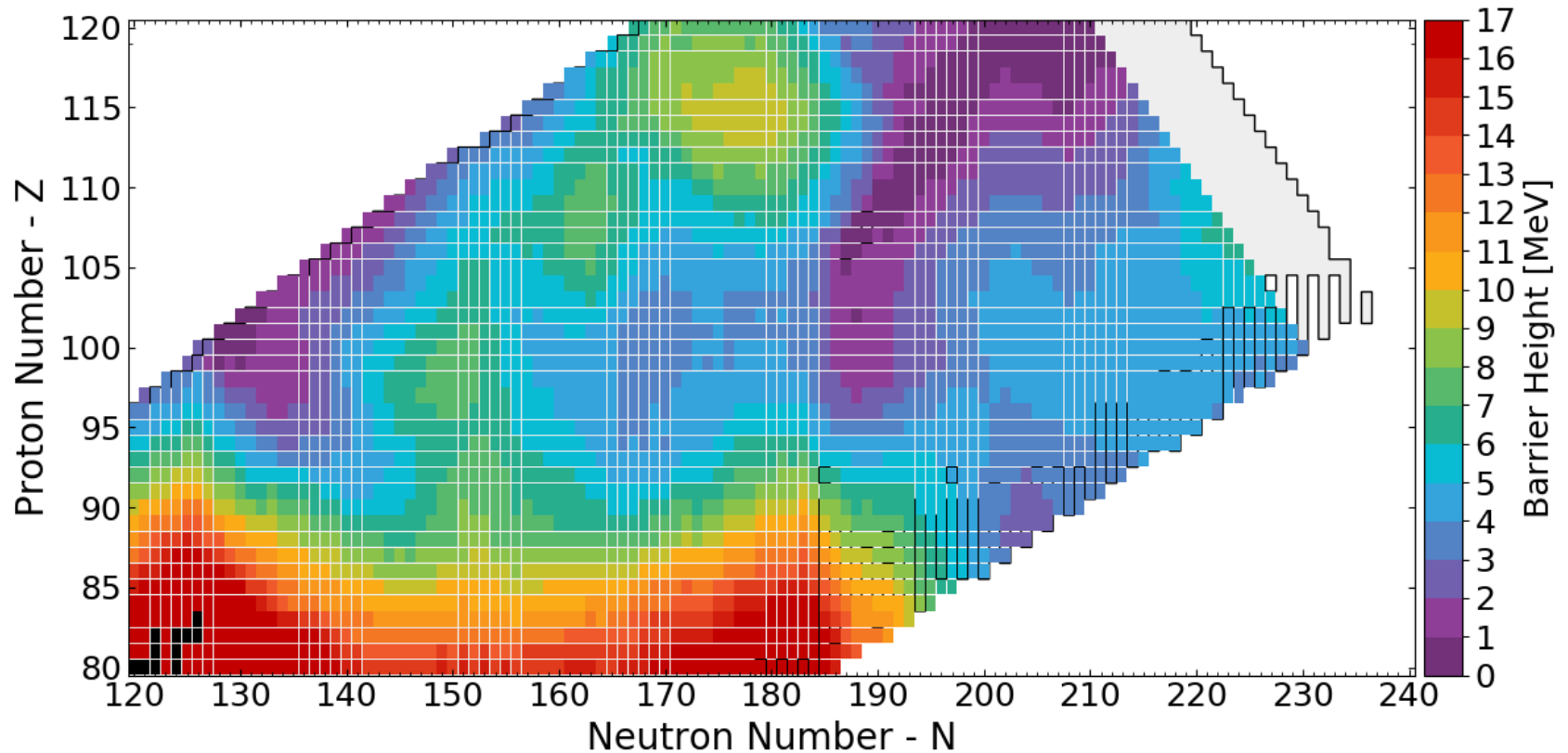
Large **Q-value**  $\Rightarrow$  impacts thermalization and therefore possibly **observations**

Responsible for what is left in the heavy mass region when nucleosynthesis is complete  $\Rightarrow$  "**smoking gun**"



# **FISSION BARRIERS**

# FISSION BARRIER HEIGHTS (FRLDM)

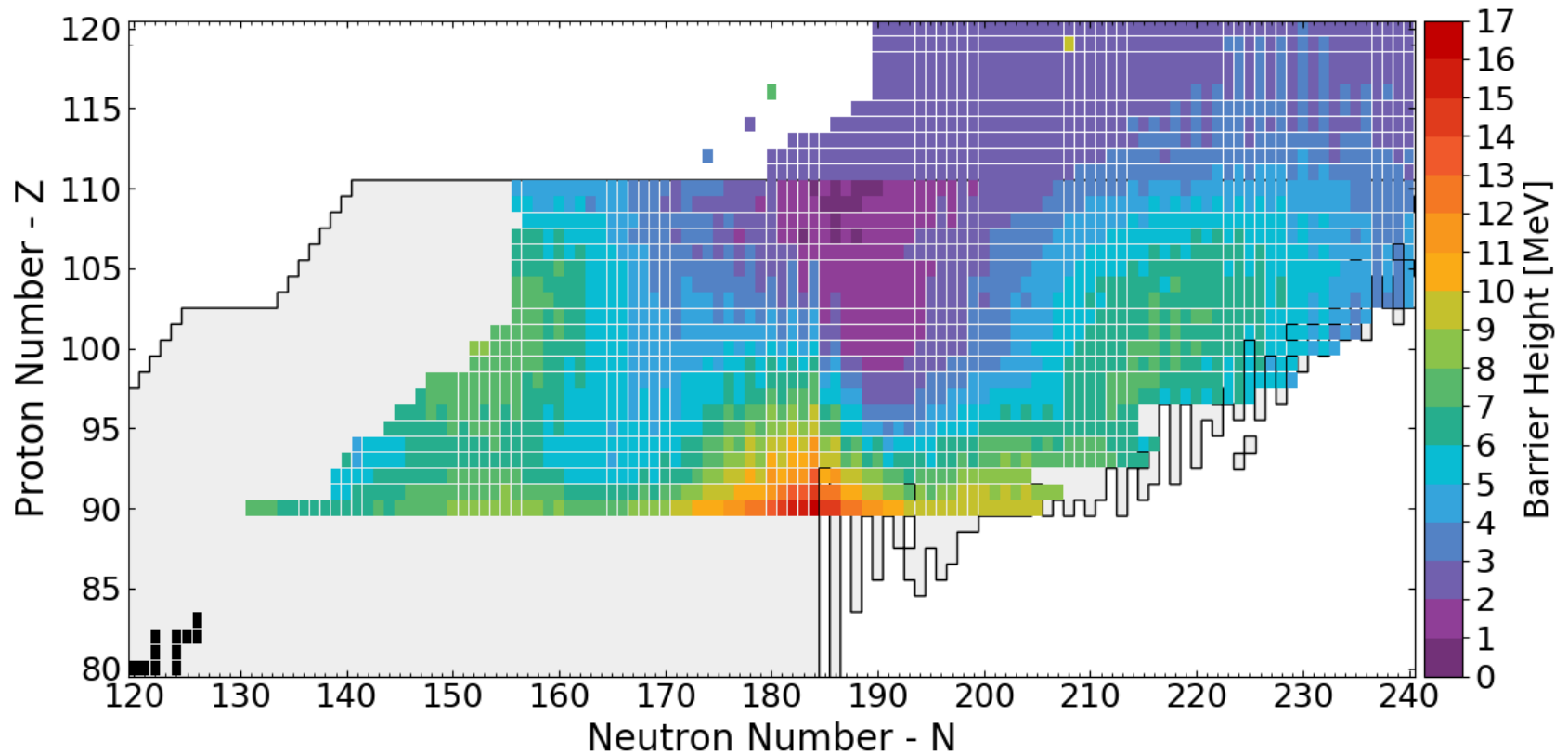


(Maximum) FRLDM Barrier heights

Low barrier  $\rightarrow$  fission  $\uparrow$

*r*-process hot spots follow low barriers

# FISSION BARRIER HEIGHTS (HFB-14)



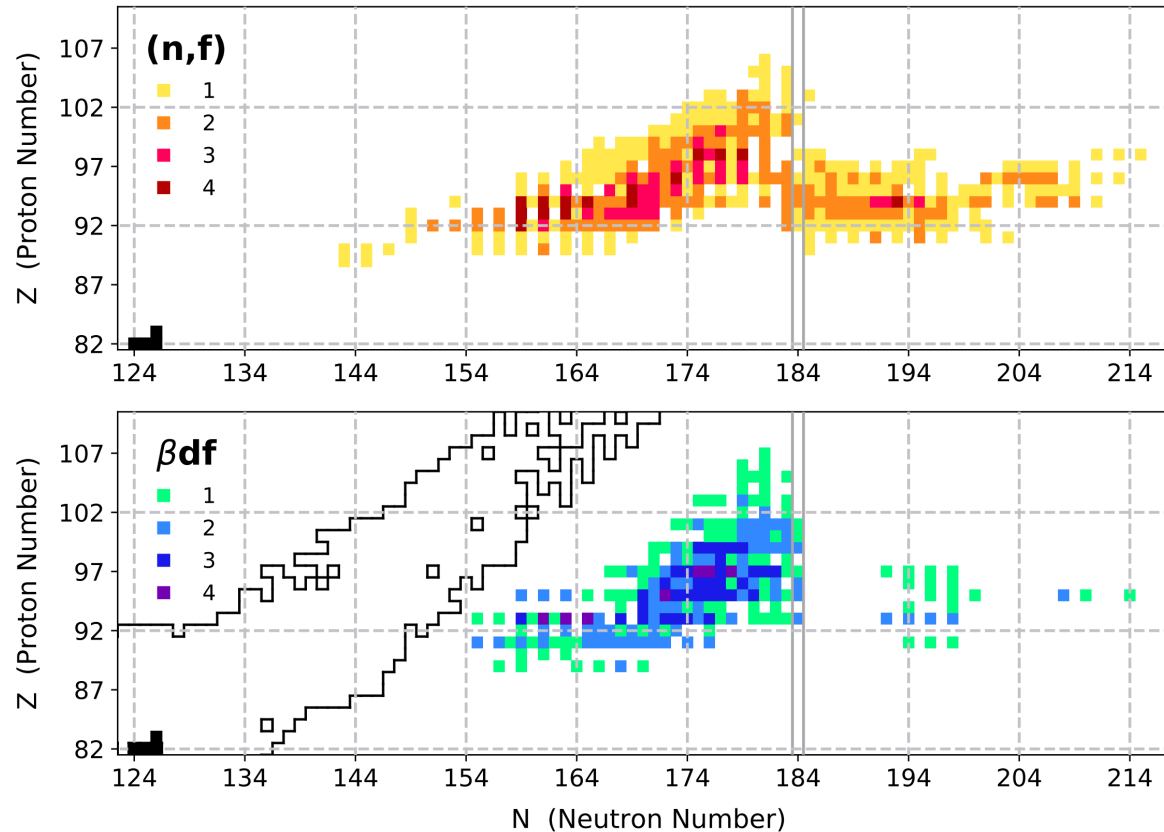
(Maximum) HFB-14 Barrier heights

Low barrier  $\rightarrow$  fission  $\uparrow$

Used in (n,f) and  $\beta$ df calculations



# FISSION HOT SPOTS



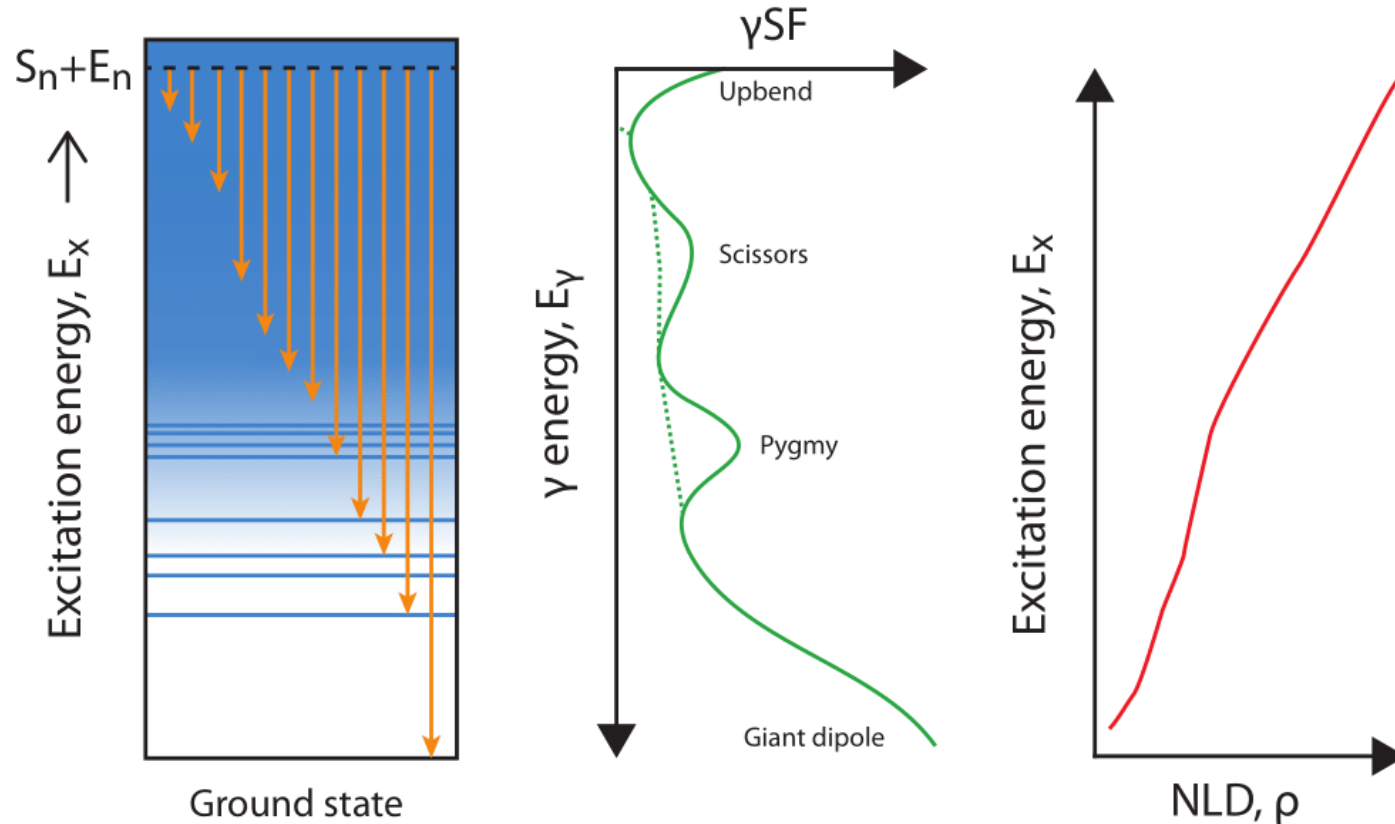
We've taken a look at the region where fission seems to occur the most

With variations in both astrophysical conditions and nuclear models

Nuclei which influence the final abundances are colored for  $(n,f)$  and  $(\beta,f)$

# NEUTRON-INDUCED FISSION

# NEUTRON-INDUCED FISSION THEORY

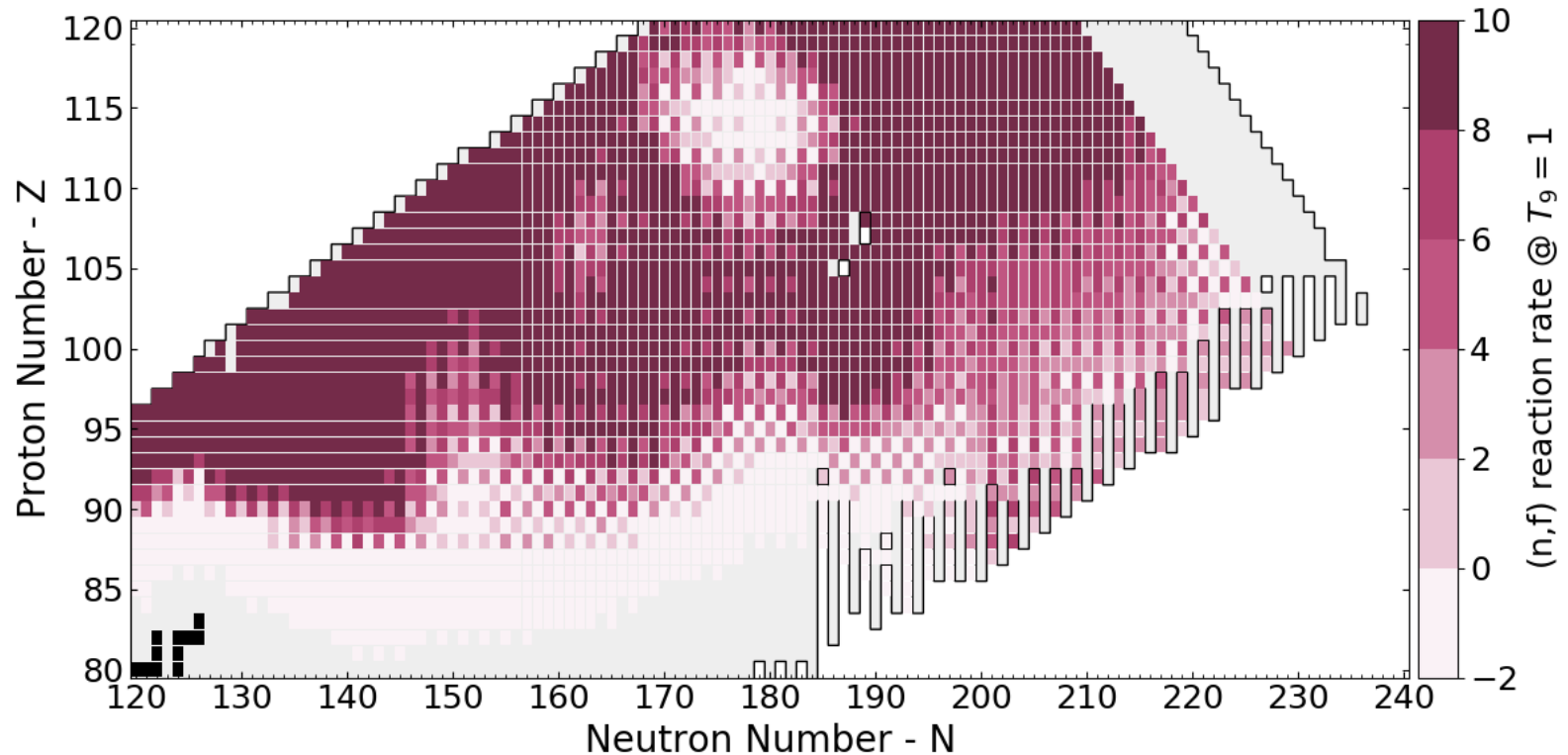


We use statistical Hauser-Feshbach theory updated to account for fission transmission

**Model inputs:** nuclear level density,  $\gamma$ -ray strength functions, optical model potentials and fission barriers



# NEUTRON-INDUCED FISSION

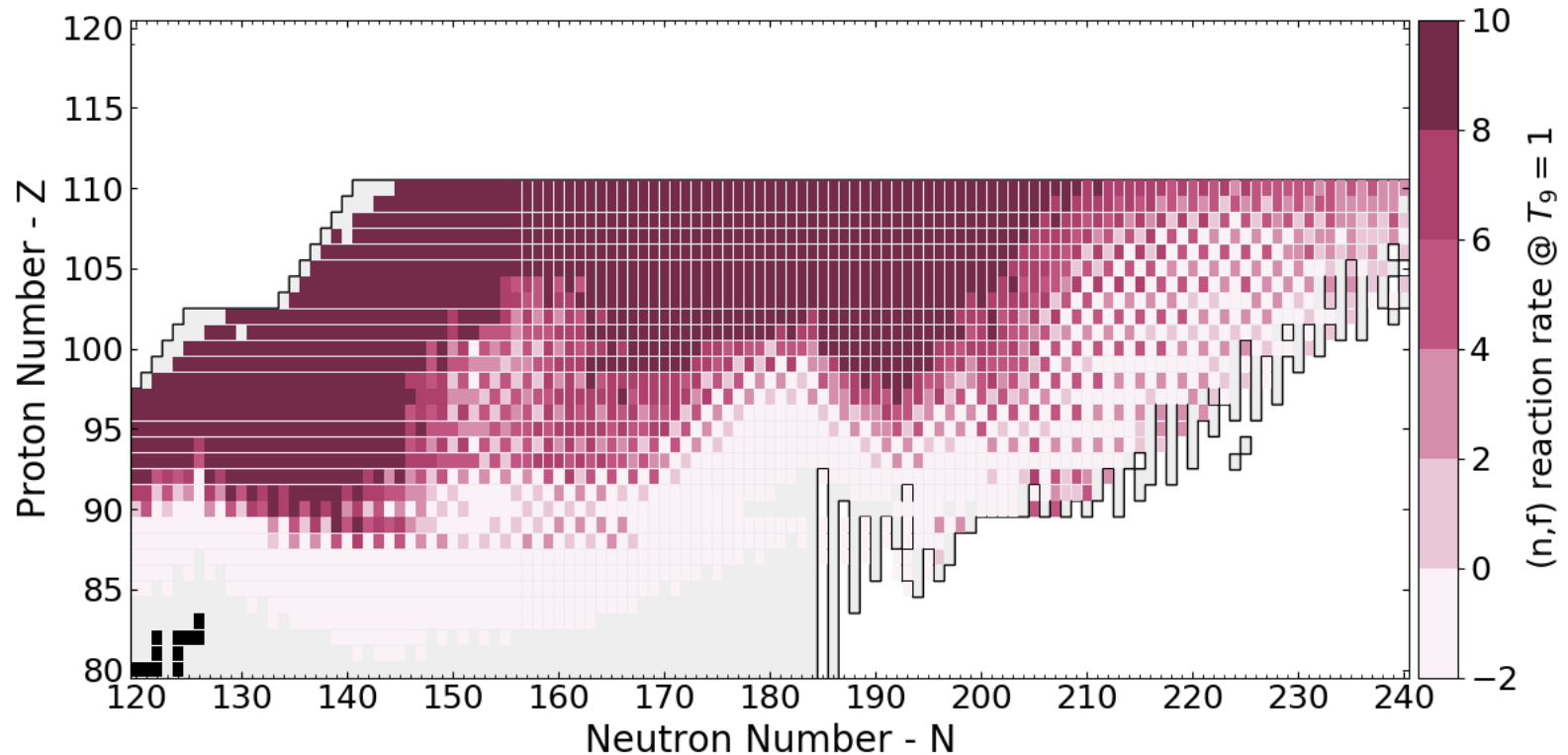


Use statistical Hauser-Feshbach for competition between neutrons,  $\gamma$ s and fission

Barrier heights Möller (2015) / FRDM2012 masses

Large region that will fission cycle  $r$ -process

# NEUTRON-INDUCED FISSION



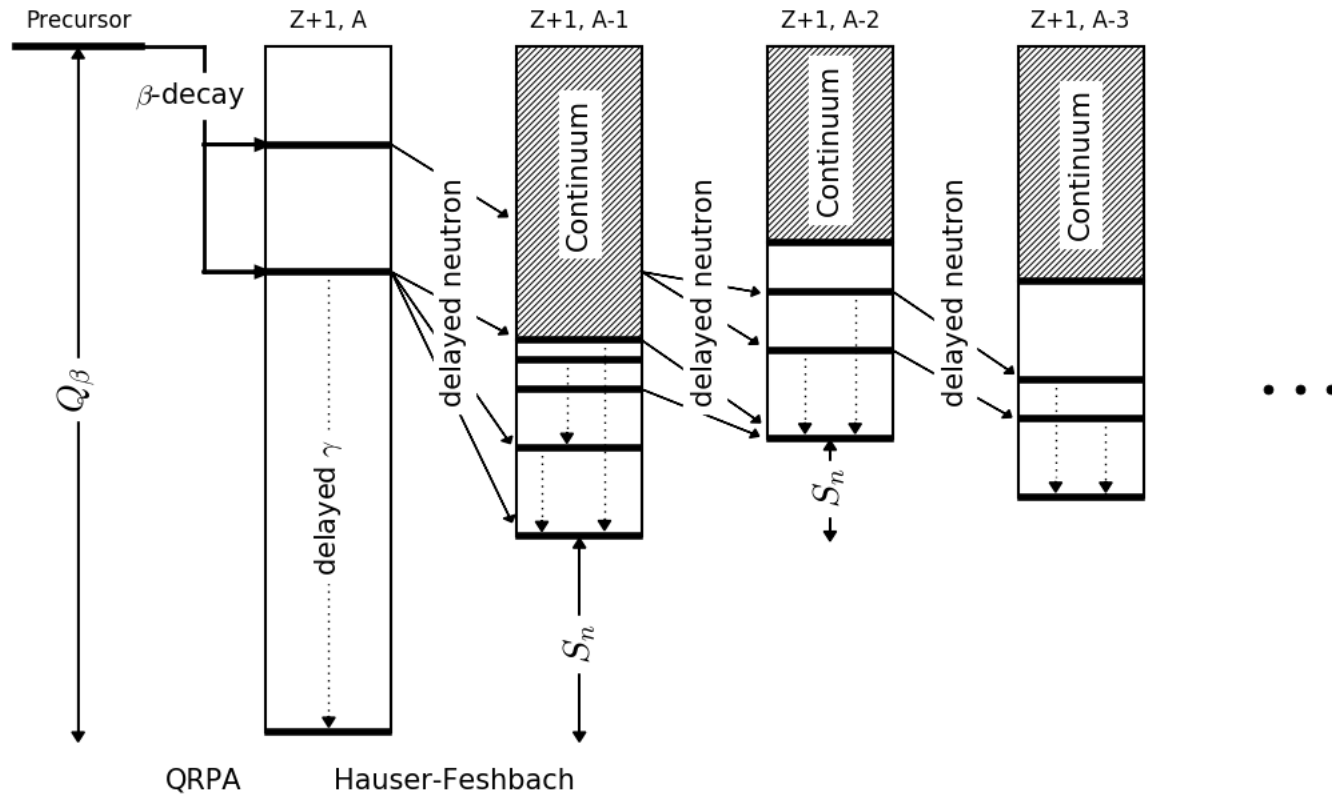
Use statistical Hauser-Feshbach for competition between neutrons,  $\gamma$ s and fission

Barrier heights from HFB-14 / HFB-17 masses

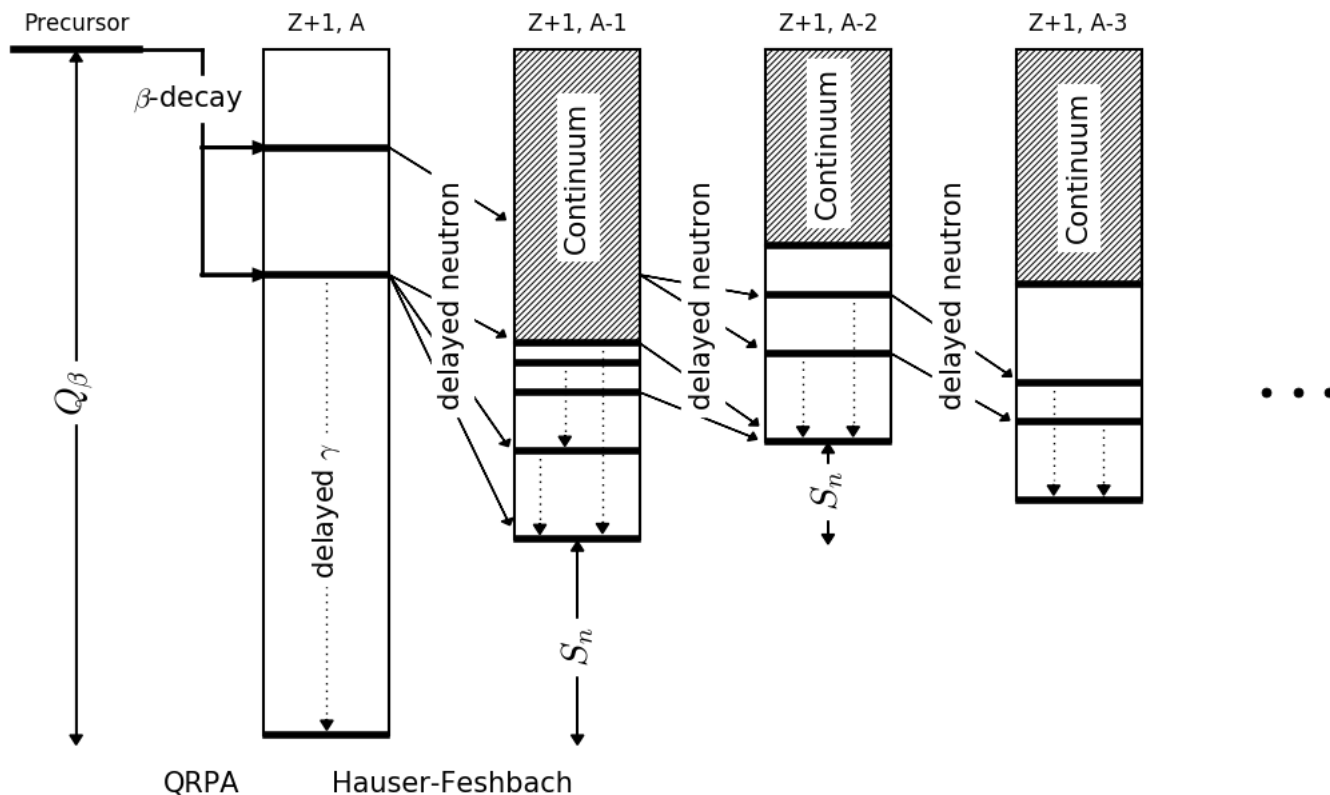
Similar results are obtained for other nuclear models

## **$\beta$ -DECAY & $\beta$ -DELAYED FISSION**

# COMBINING QRPA + HF



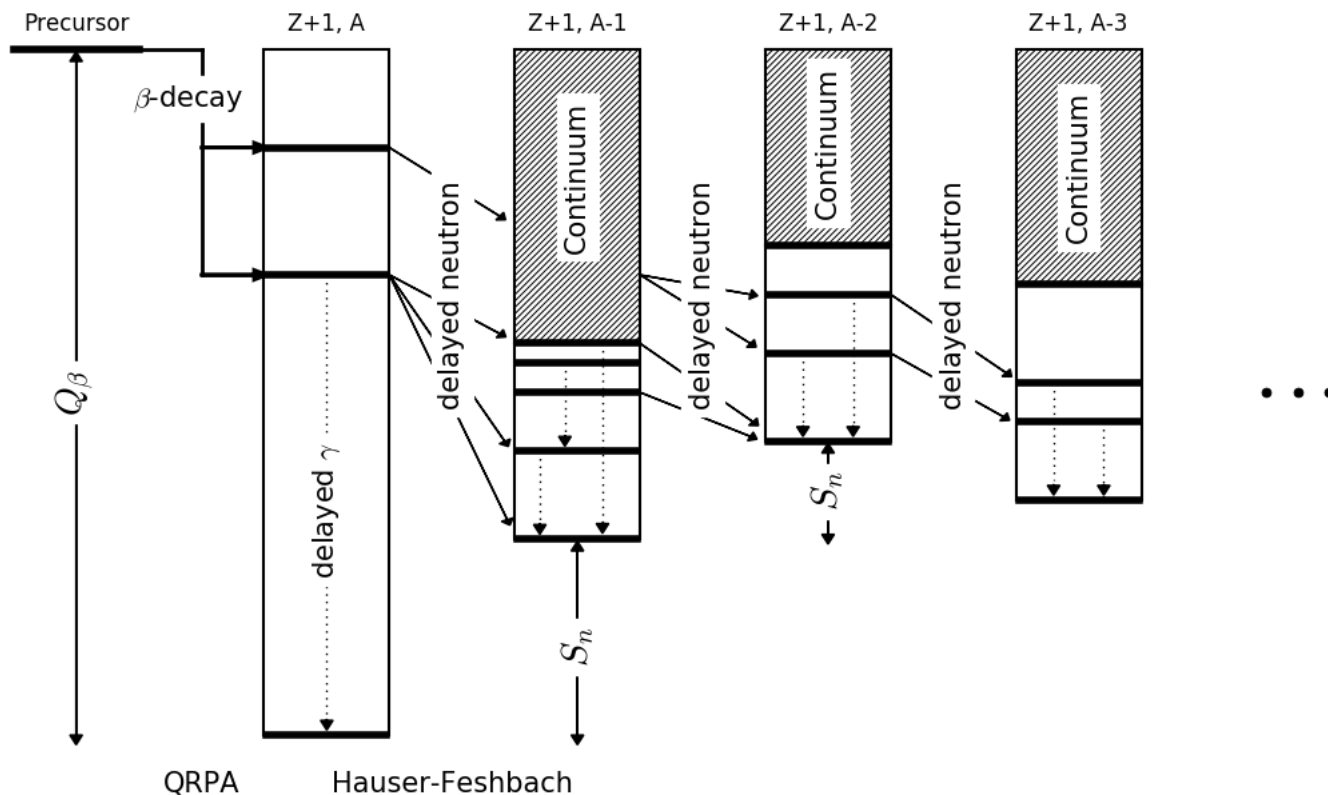
# COMBINING QRPA + HF



Initial population from the  $\beta$ -decay strength function from P. Möller's QRPA



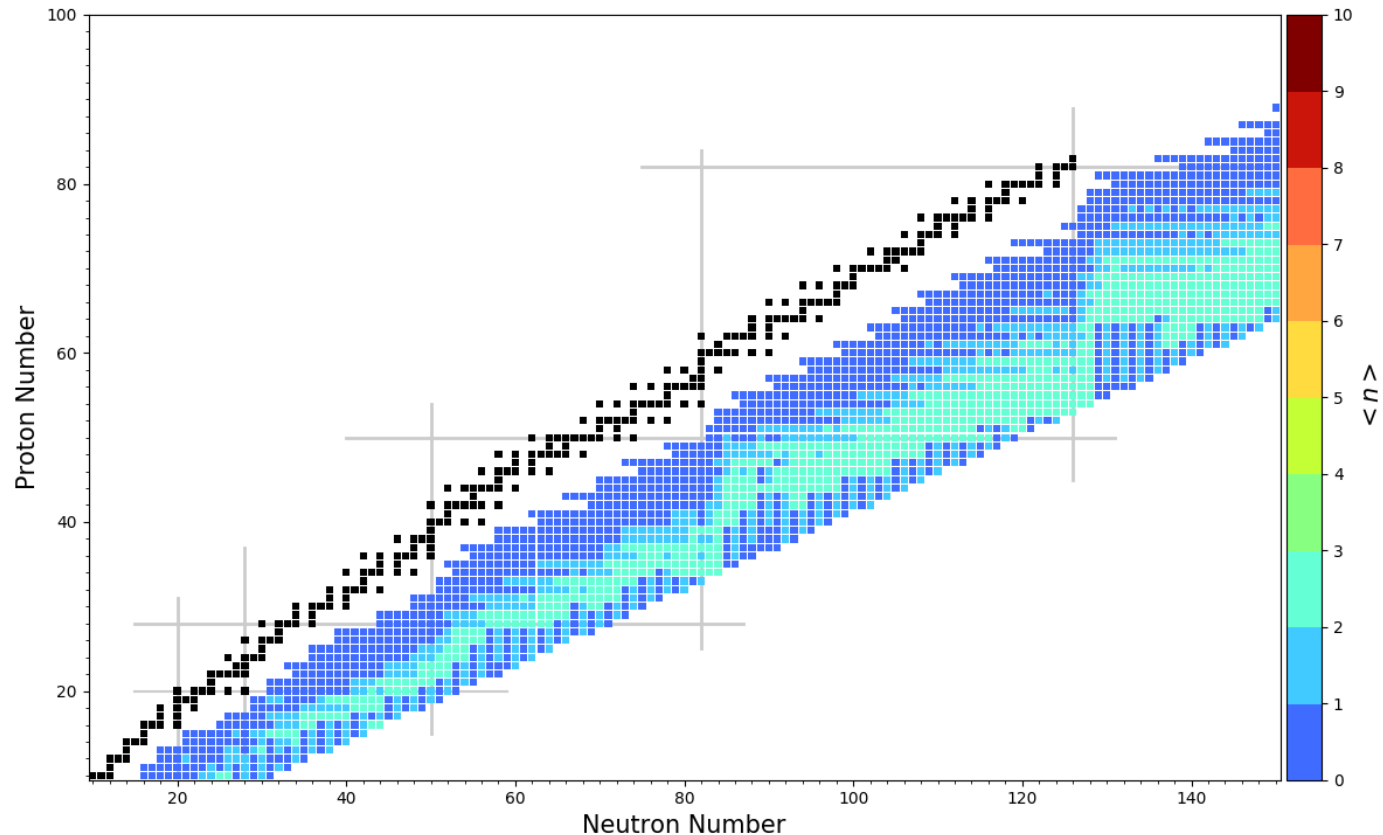
# COMBINING QRPA + HF



Initial population from the  $\beta$ -decay strength function from P. Möller's QRPA

Follow the statistical decay until all excitation energy is exhausted

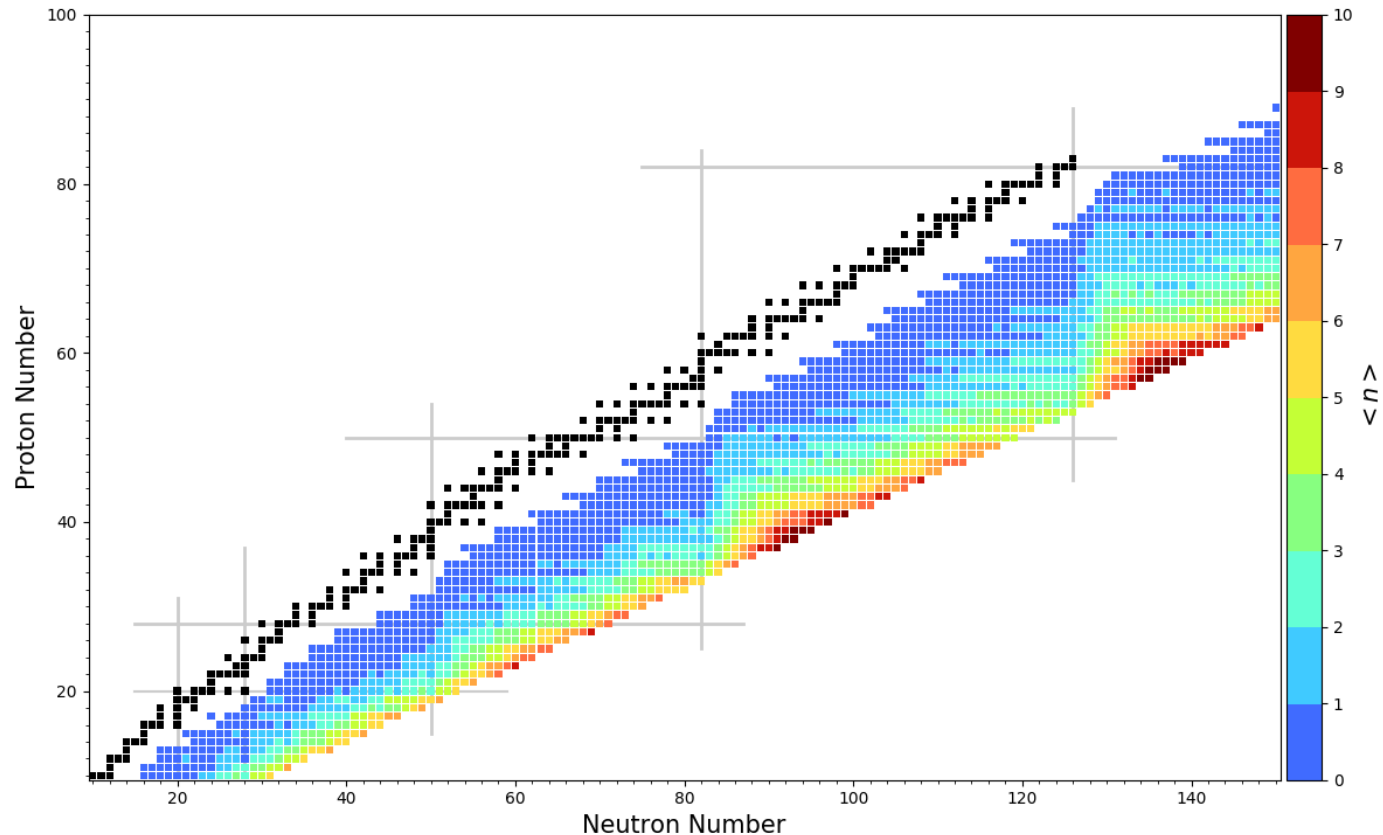
# AVERAGE NEUTRON EMISSION



Apply energy window method to the entire chart of nuclides

Problem with describing very neutron-rich nuclei

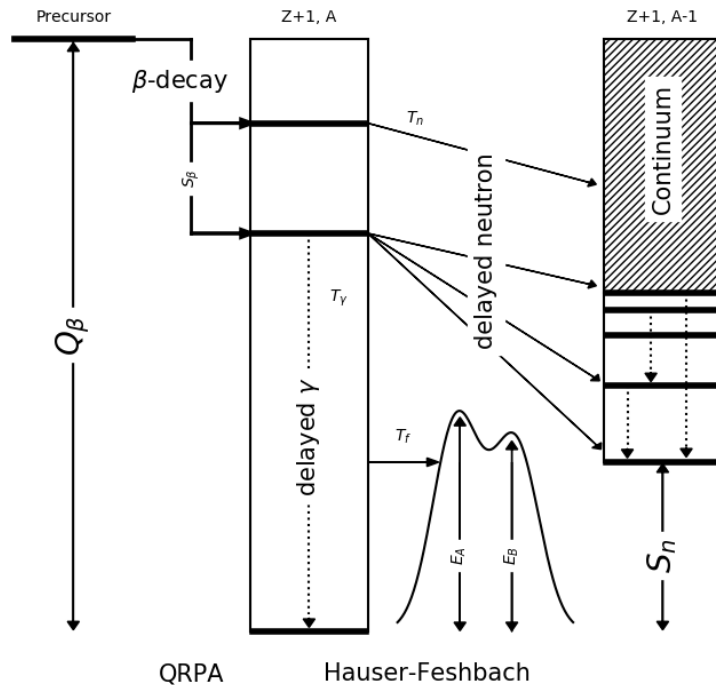
# AVERAGE NEUTRON EMISSION



Apply the [QRPA+HF](#) method to the entire chart of nuclides

Problem with neutron-rich nuclei goes away

# $\beta$ -DELAYED FISSION



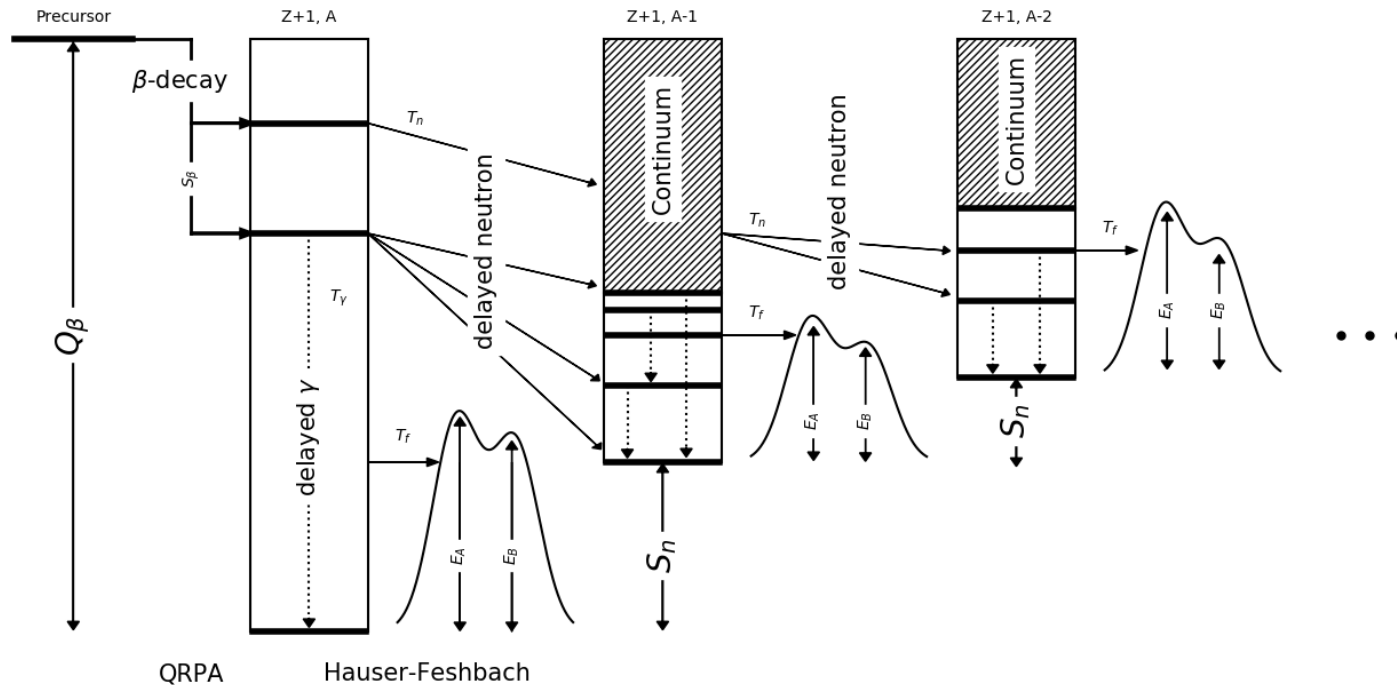
We have recently extended our QRPA+HF model to describe  $\beta$ -delayed fission ( $\beta$ df)

Barrier heights from Möller *et al.* PRC 91 024310 (2015)

Assumes a Hill-Wheeler form for fission transmission

Mumpower *et al.* PRC 94 064317 (2016) • Spyrou *et al.* PRL (2016) • Möller *et al.* ADNDT 125 (2019)  
Yokoyama *et al.* PRC (2019) • Mumpower *et al.* ApJ (2018)

# MULTI-CHANCE $\beta$ DF

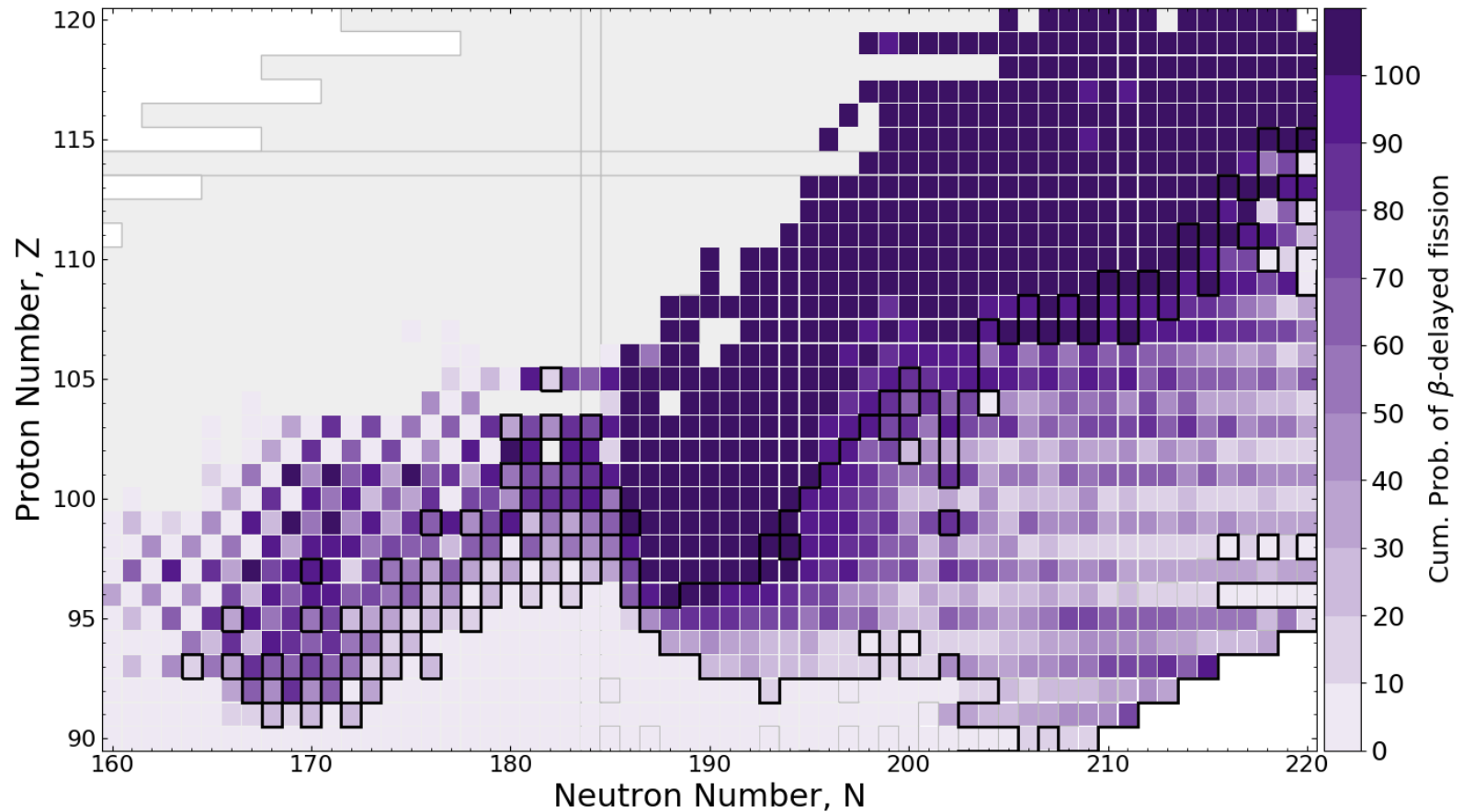


Recall: Near the dripline  $Q_{\beta} \uparrow$   $S_n \downarrow$

Multi-chance  $\beta$ df: each daughter may fission

The yields in this decay mode are a convolution of many fission yields!

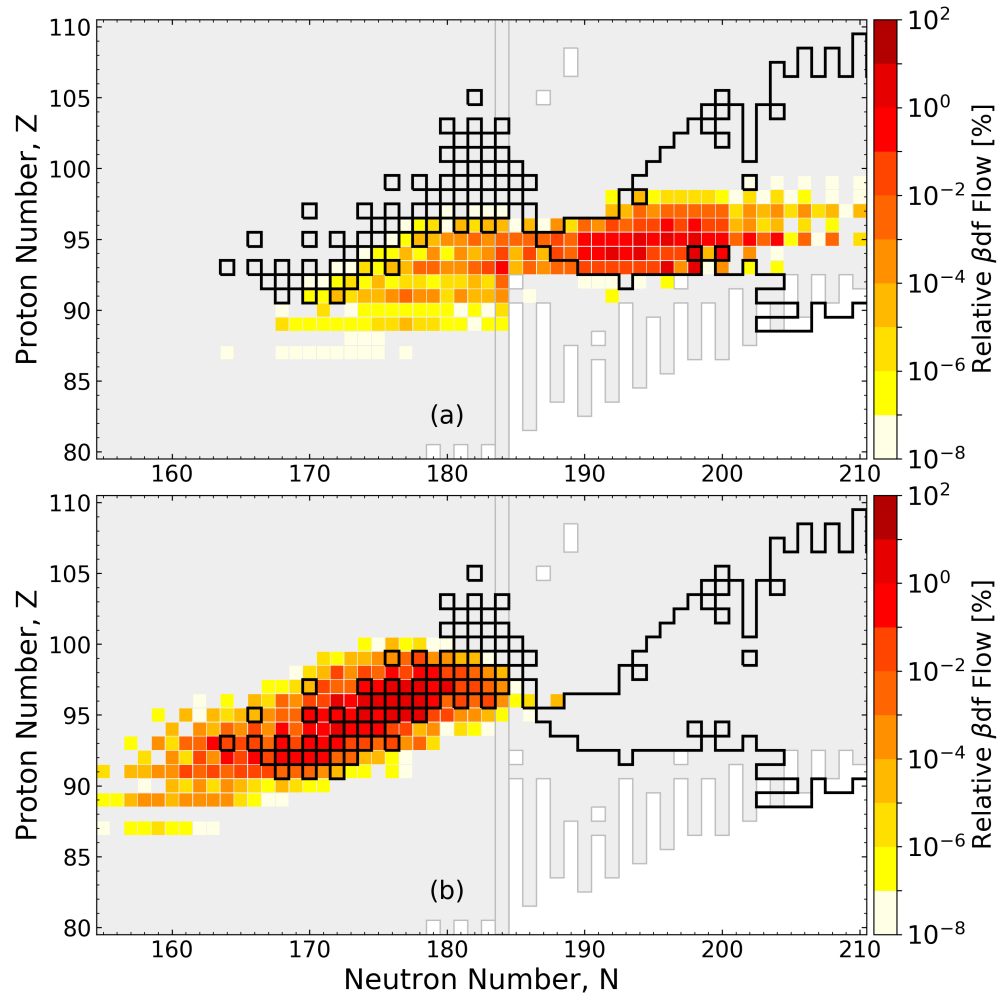
# CUMULATIVE $\beta$ DF PROBABILITY



$\beta$ df occupies a large amount of real estate in the NZ-plane

Multi-chance  $\beta$ df outlined in black

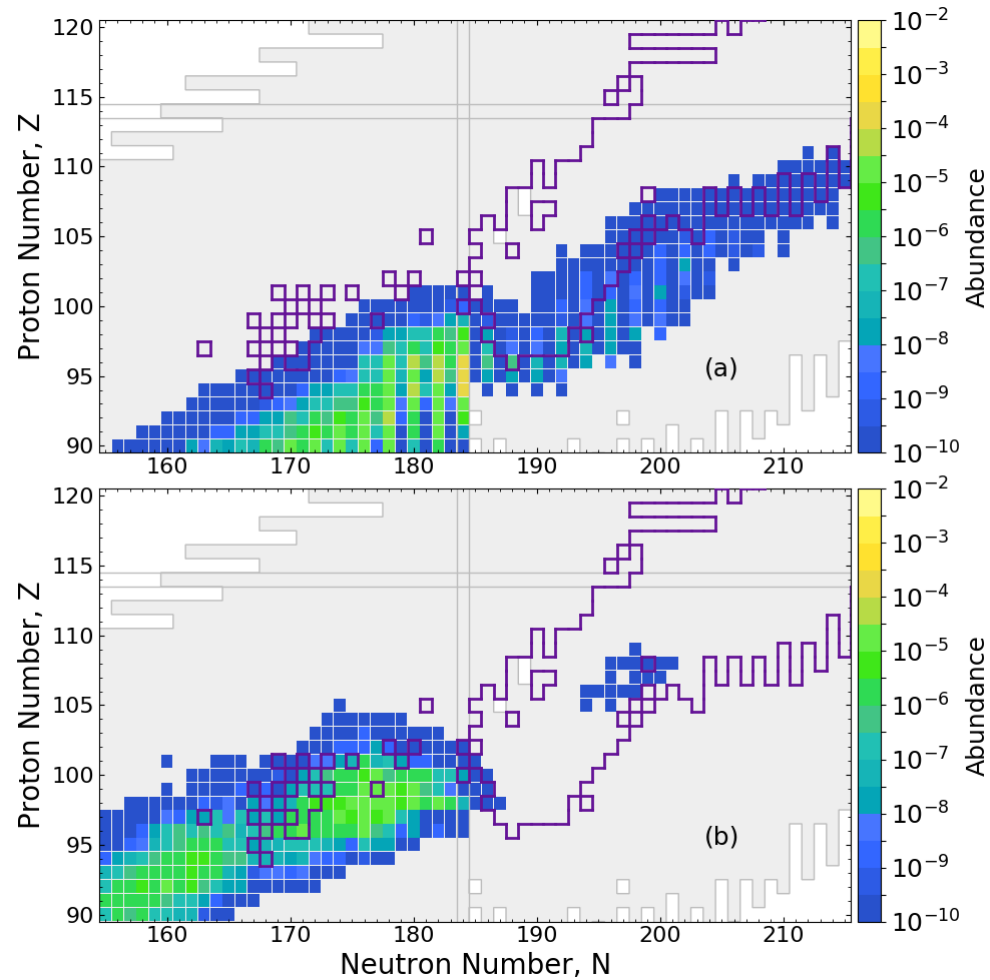
# MULTI-CHANCE $\beta$ DF CONTRIBUTION



Network calculation of neutron star merger ejecta; FRDM2012 inputs

Multi-chance  $\beta$ df contributes at both early and late times

# APPLICATION TO THE $r$ -PROCESS

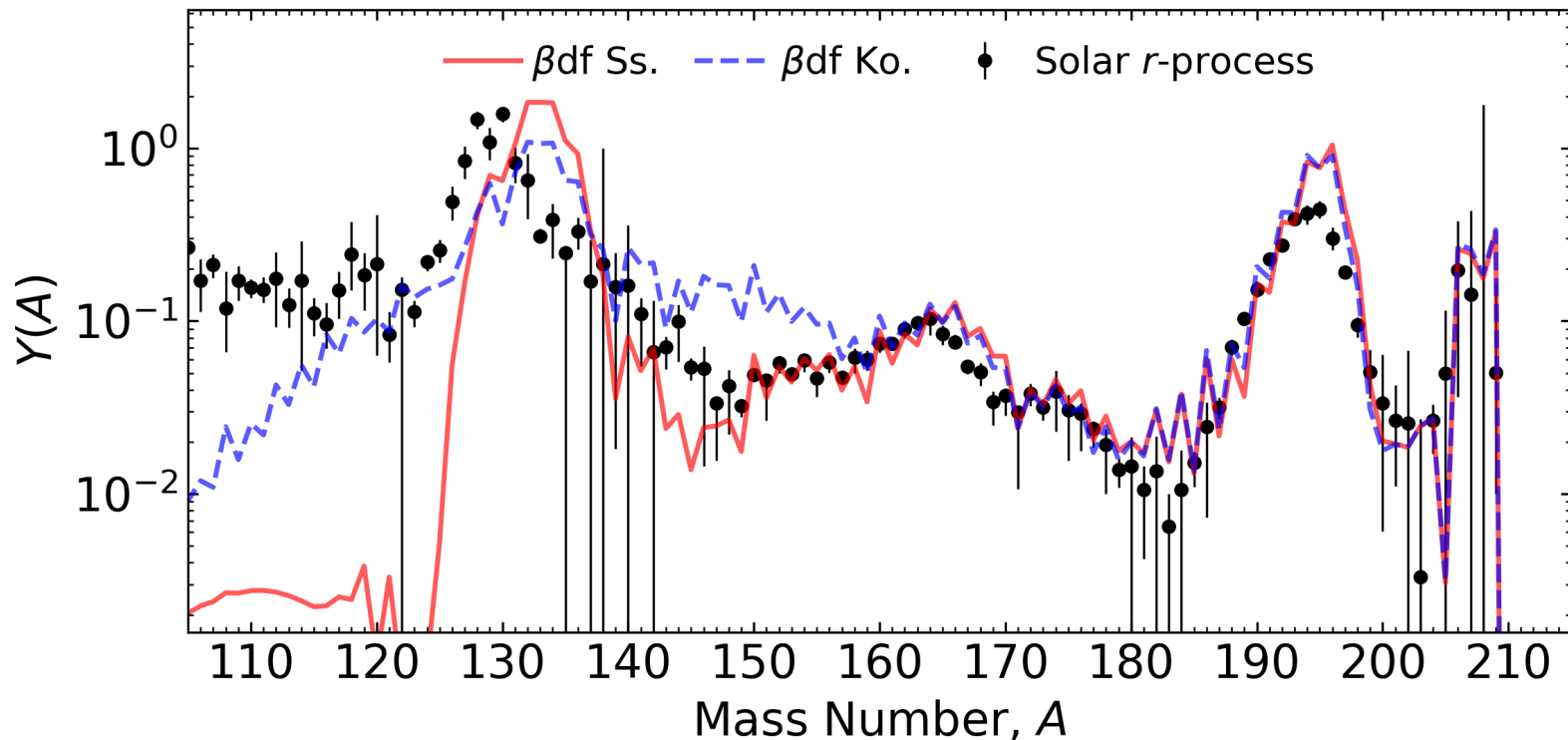


PRISM network calculation of neutron star merger ejecta

$\beta$ df alone may prevent the production of superheavy elements in nature



# IMPACT ON FINAL ABUNDANCES



Network calculation of tidal ejecta from a neutron star merger (FRDM2012)

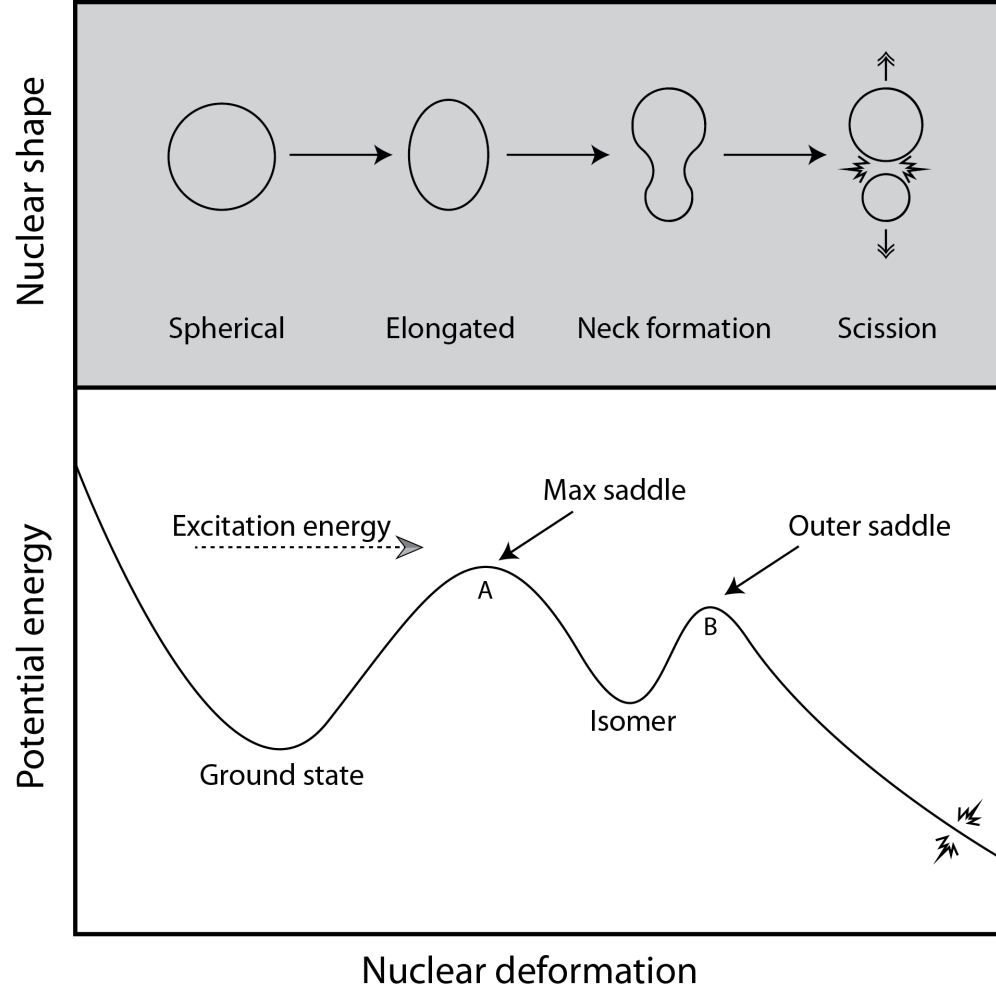
$\beta df$  can shape the final pattern near the  $A = 130$  peak

This is because of a relatively long fission timescale

Conclusion  $\Rightarrow$  we need a good description of fission yields to understand abundances near  $A \sim 130$ .

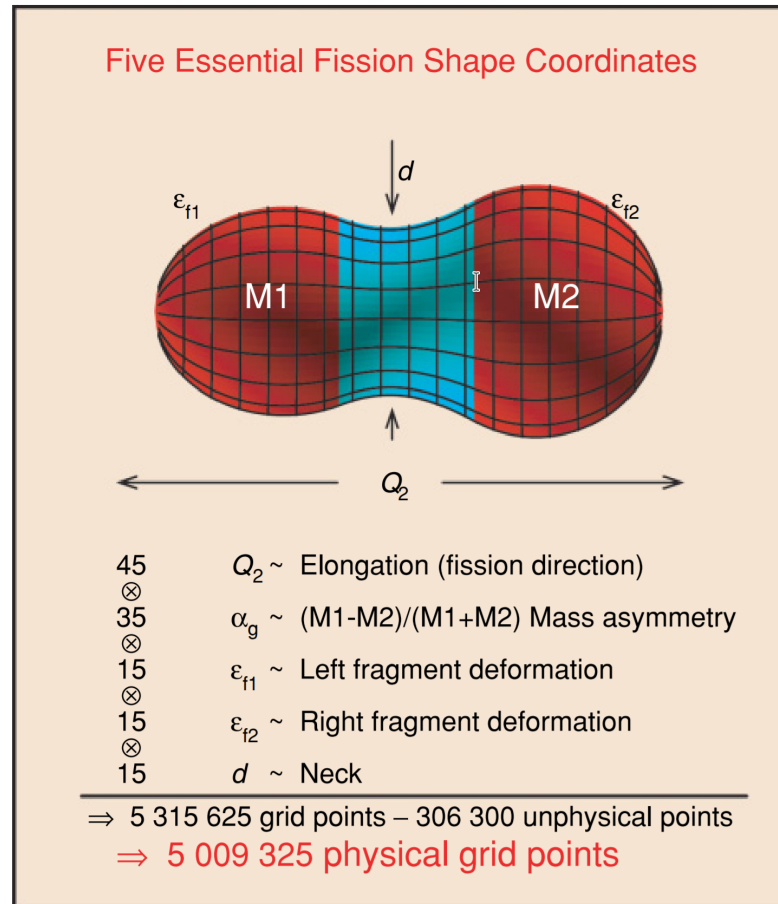
# FISSION YIELDS

# A SIMPLE PICTURE OF FISSION



Follow progression of the nucleus from compact to highly elongated shapes

# FINITE-RANGE LIQUID-DROP MODEL



Many possible shape degrees of freedom - but we have to isolate the most important

# HOW DO WE CALCULATE FRAGMENT YIELDS WITH THIS MODEL?

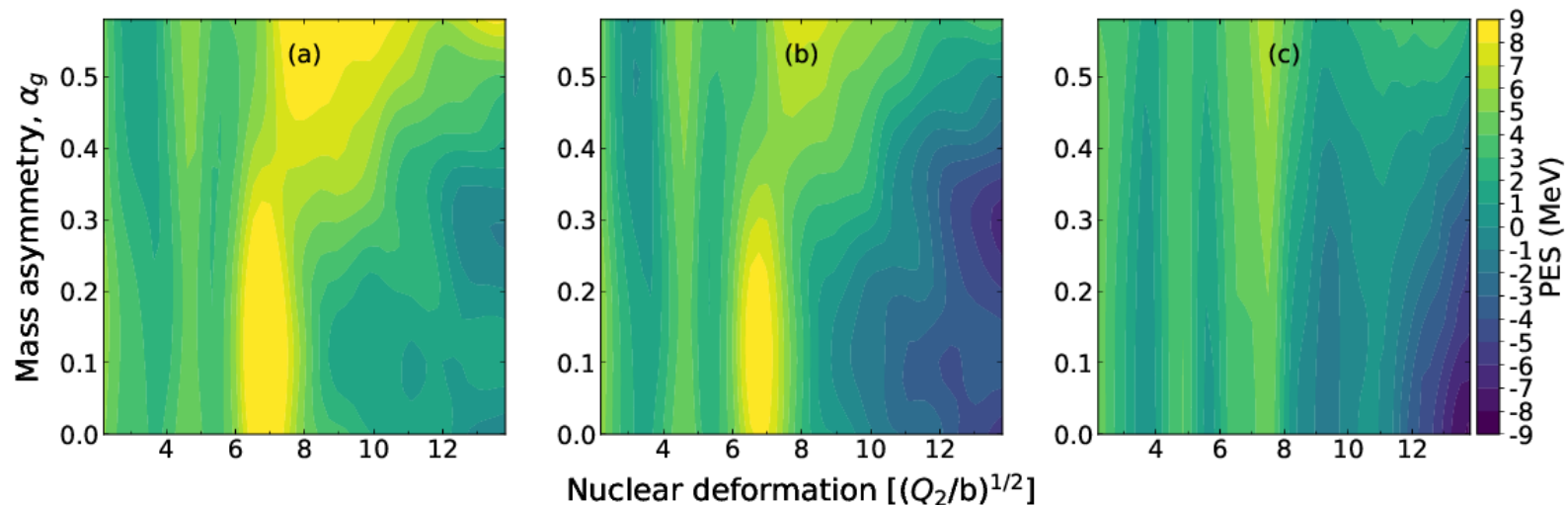
Change in nuclear shape acts as a driving force for bulk rearrangement of material

This results in a collective kinetic energy

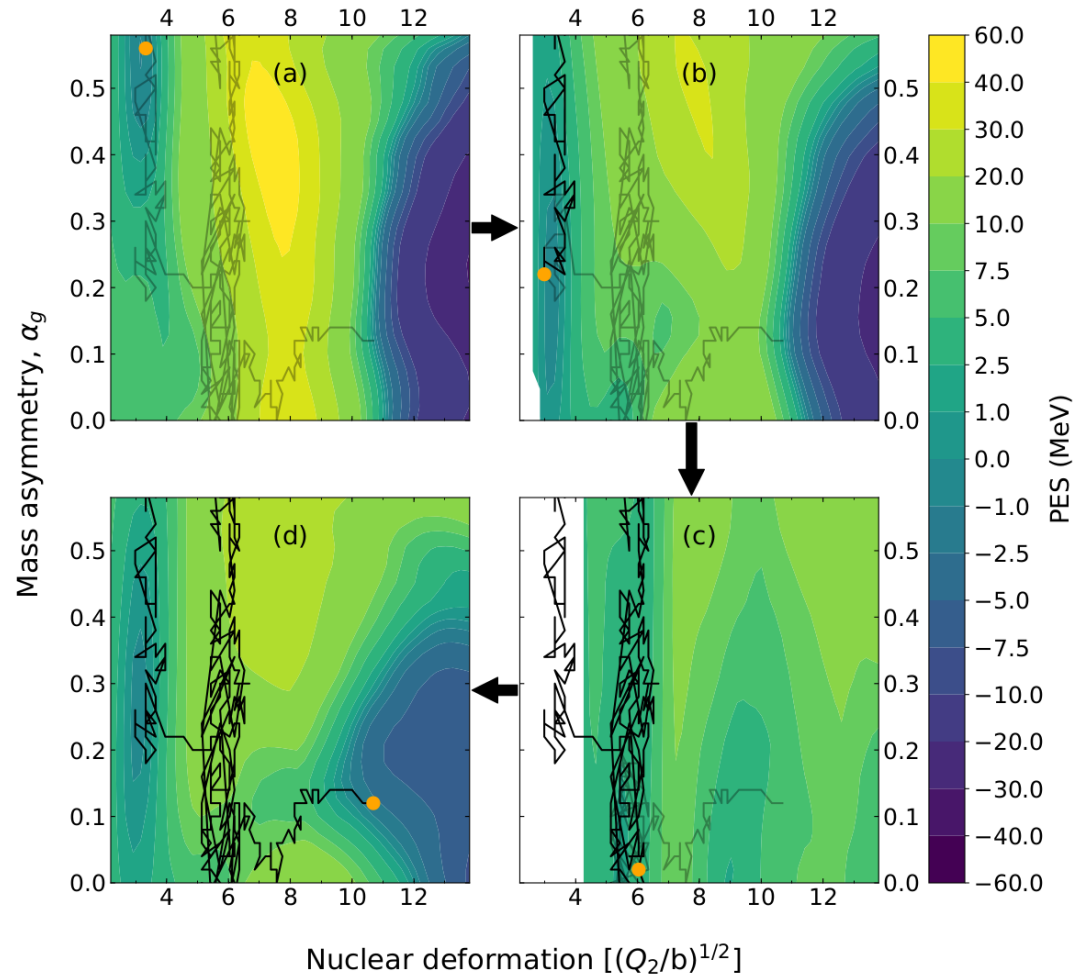
The macroscopic shape degrees of freedom couple to individual nucleonic motion

Resulting in an evolution that is both damped and diffuse

This can be approximated as **Brownian shape motion**

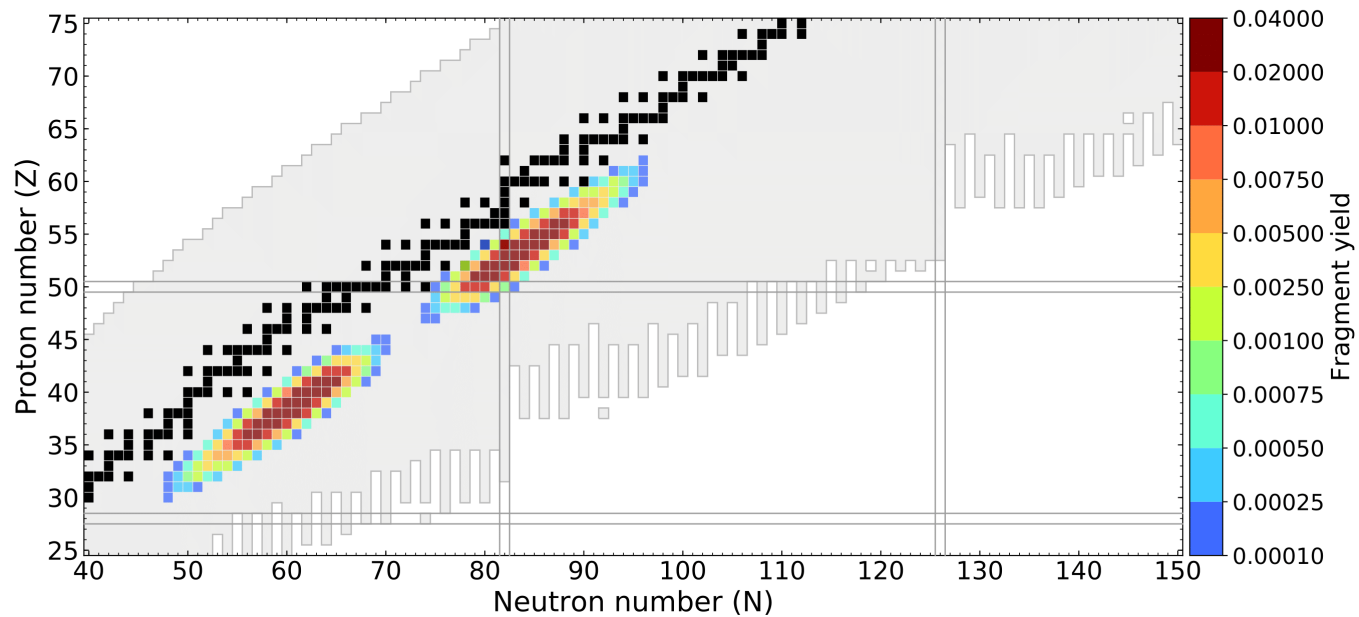
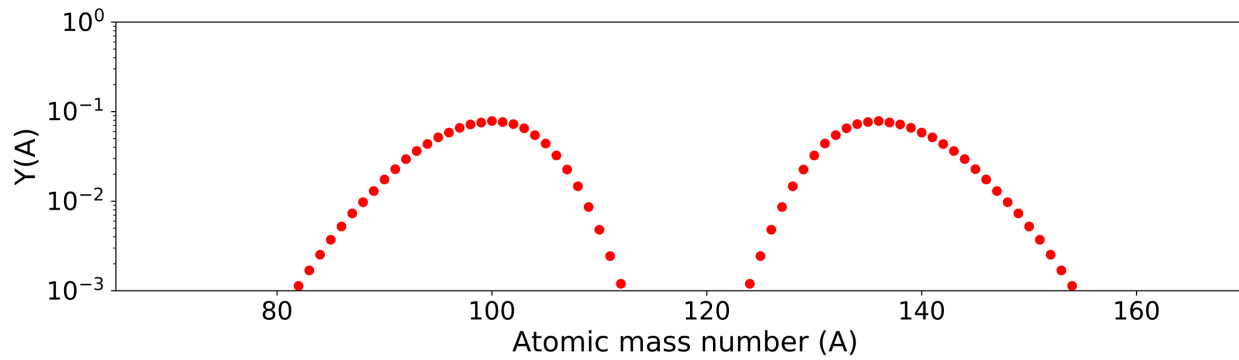


# FISSION EVOLUTION

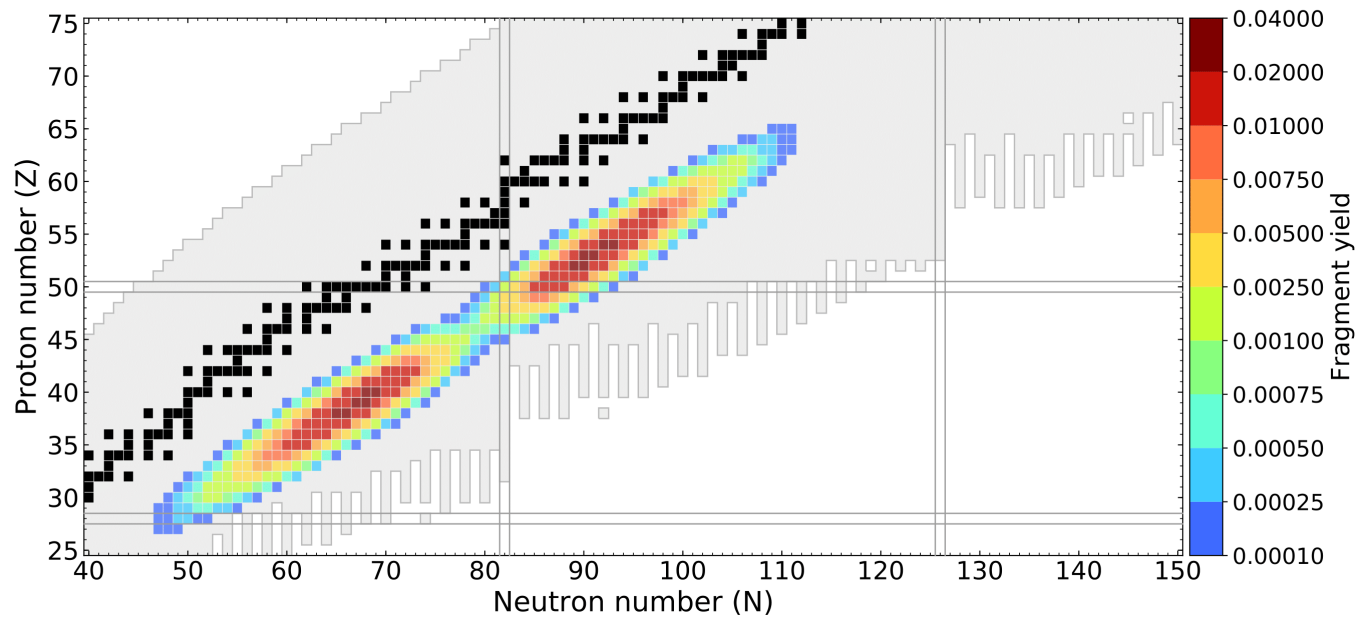
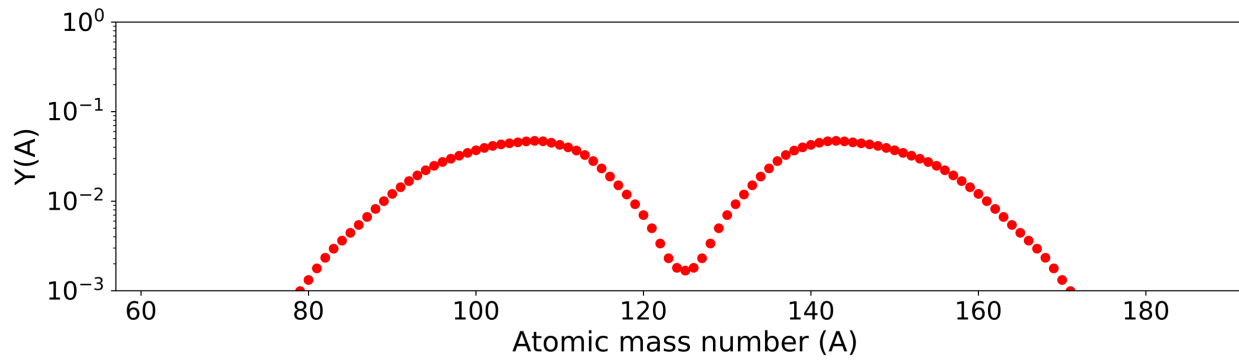


Amounts to random walk across potential energy surface

# $^{236}\text{U}$ Y(Z,A) YIELD

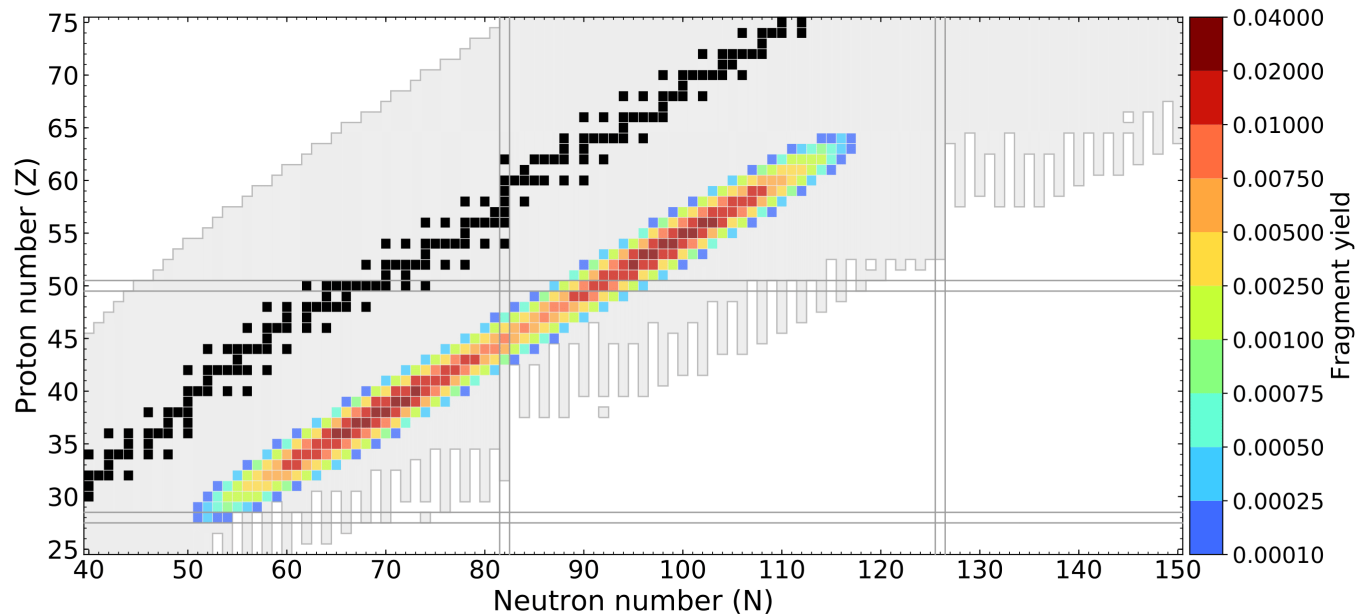
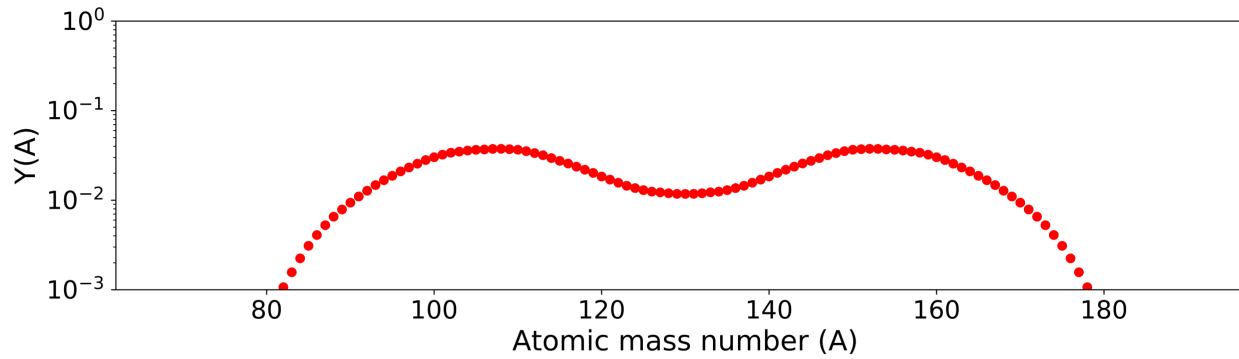


# 250-U Y(Z,A) YIELD

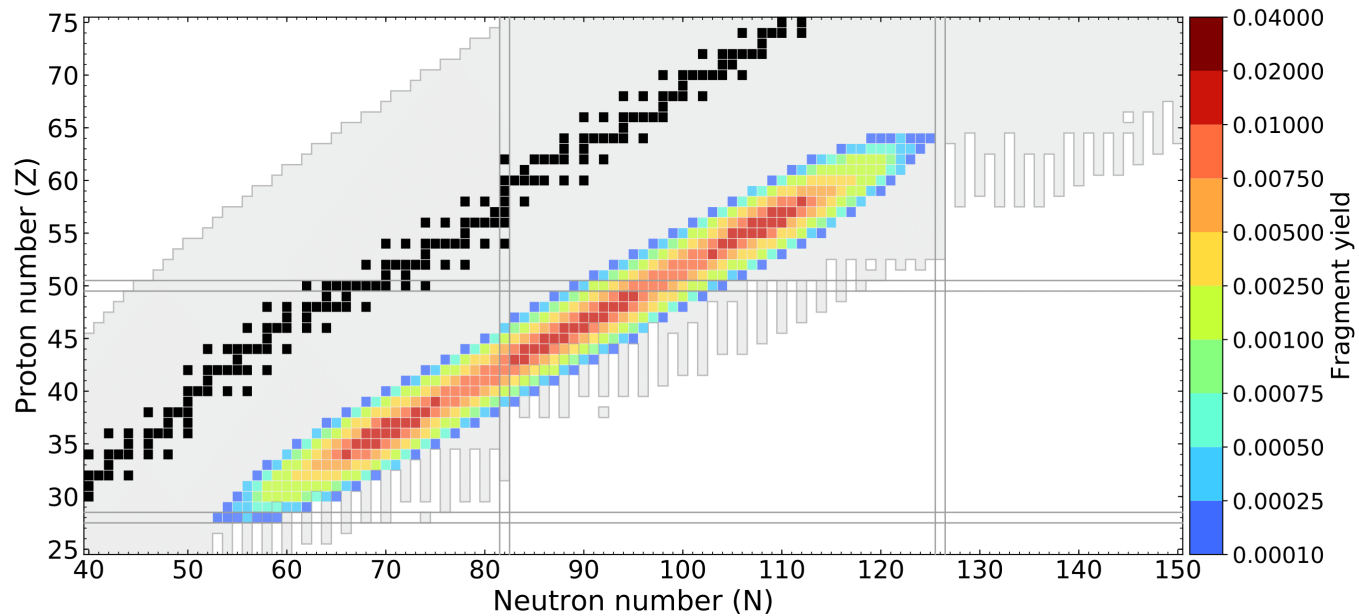
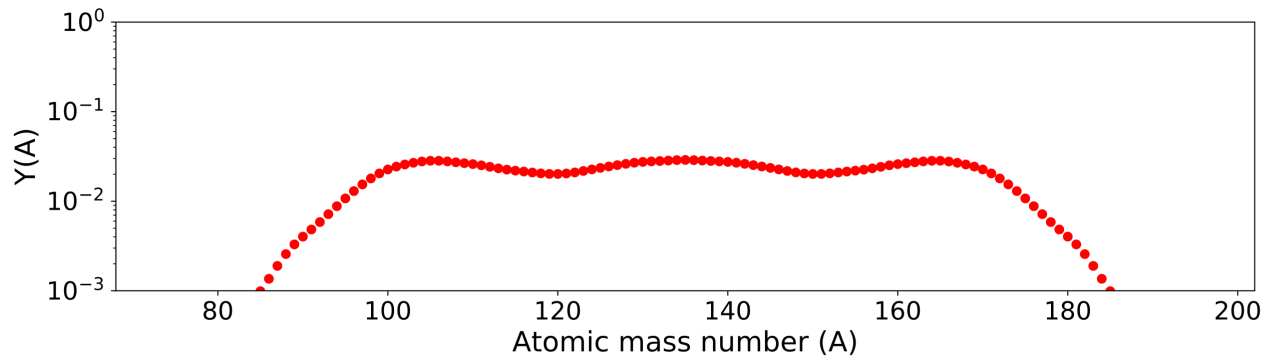




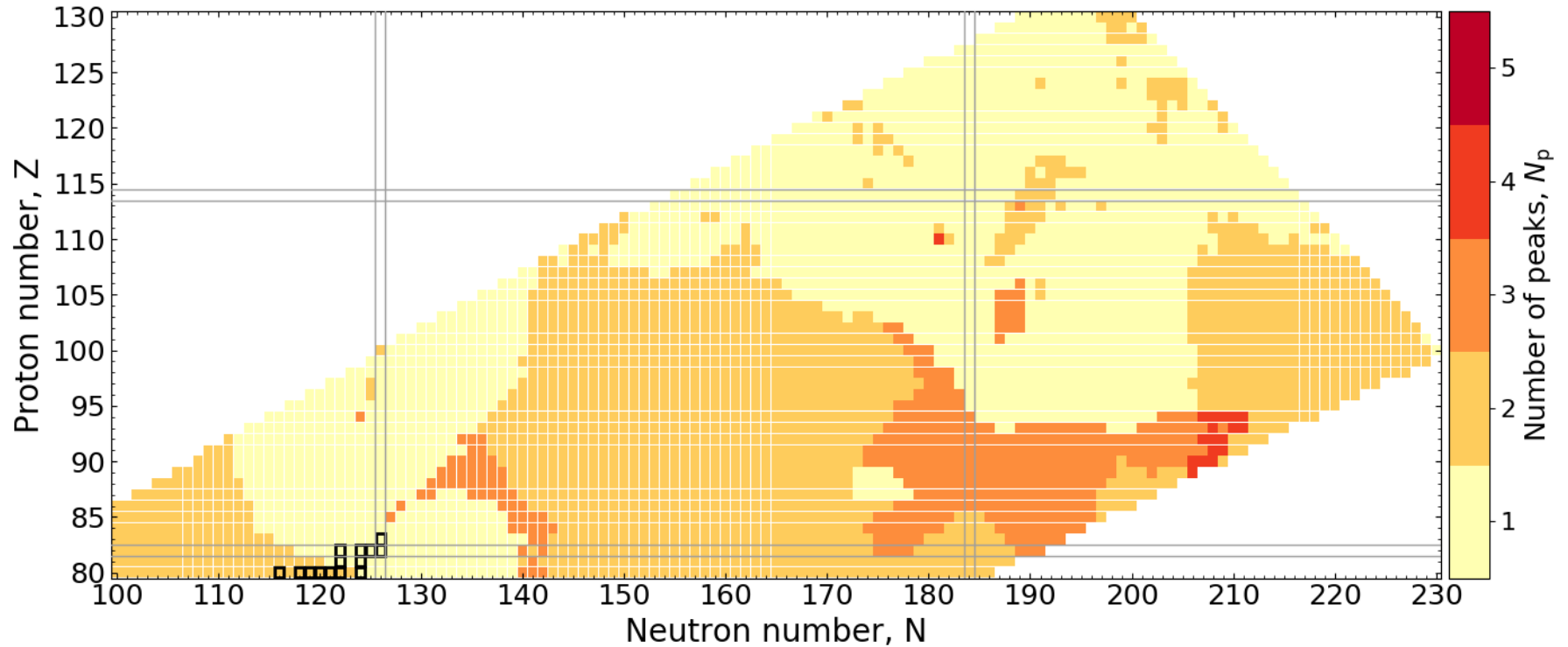
# 260-U Y(Z,A) YIELD



# 270-U Y(Z,A) YIELD



# NUMBER OF PEAKS

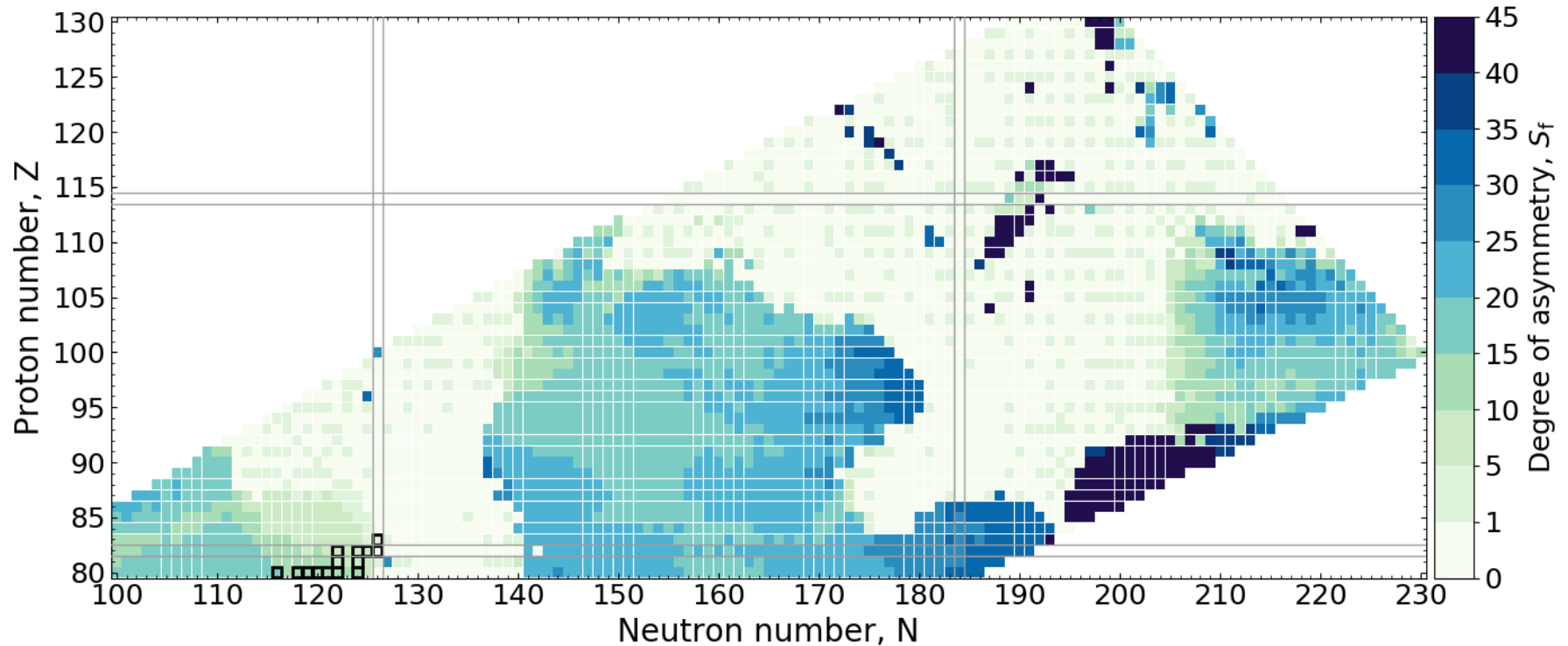


Count the number of peaks in the mass yield,  $Y(A)$ , distribution

Rather smooth variation in number of peaks across chart of nuclides.

$r$ -process region: 2 or 3 peaks are the norm given our prediction of fission hot spots

# MEASURE OF ASYMMETRY

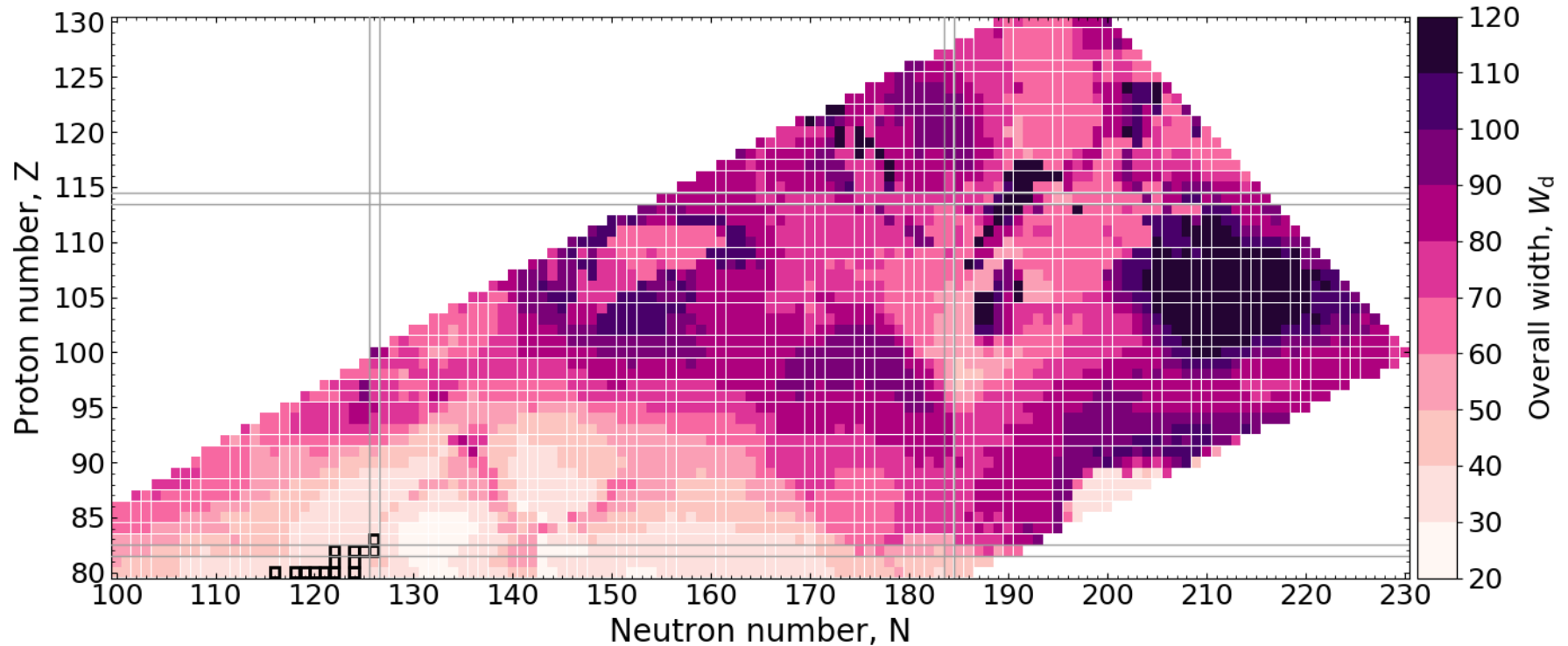


Measure the distance in  $A$  between the maxima of  $Y(A)$  and  $Y(A_f/2)$

Abrupt changes can be seen when the maxima shift from symmetric to asymmetric

Symmetric followed by asymmetric distributions can be expected in  $r$ -process simulations

# EXTENT OF Y(A) DISTRIBUTION

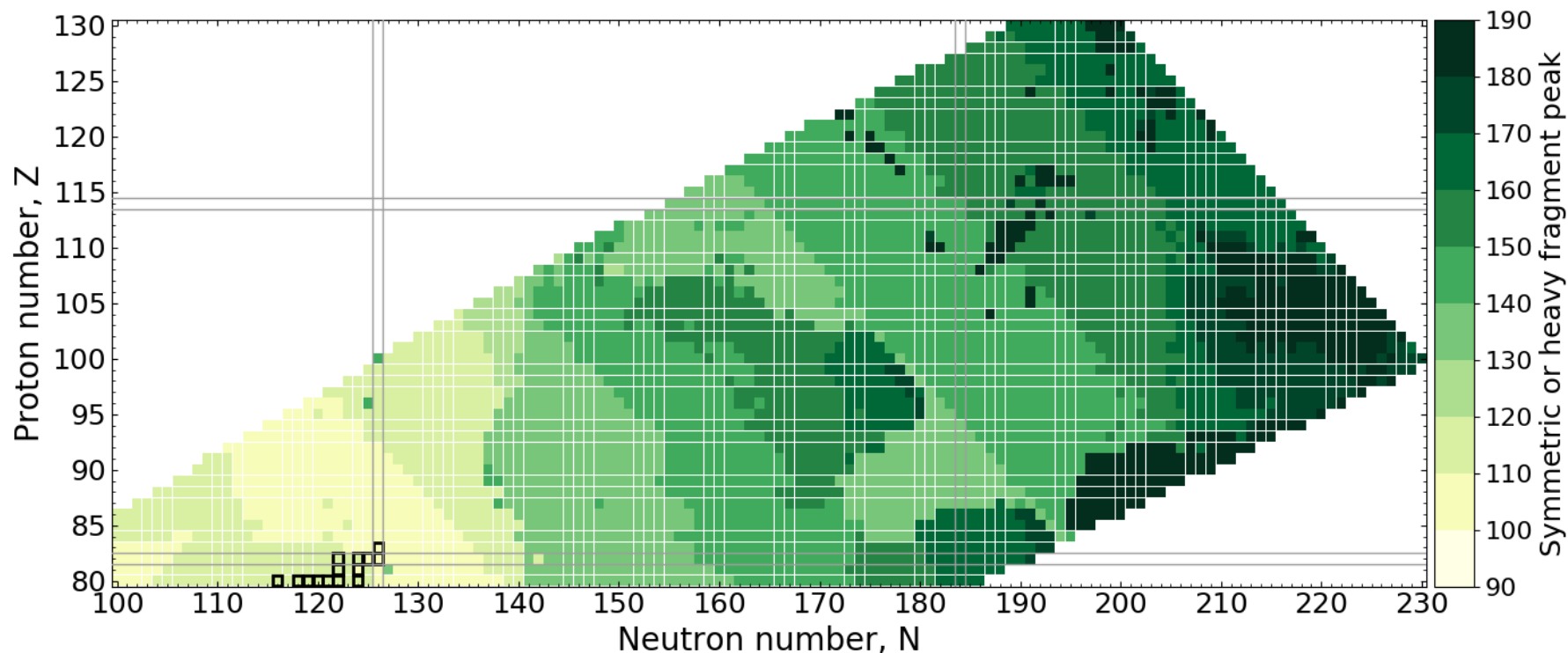


Measure the spread of the daughter products in  $A$

Strong dependence can be seen with the fission system

Wiggles in the yield (number of peaks or asymmetry) don't matter if the distribution is wide!

# PEAK LOCATION ( $A$ )

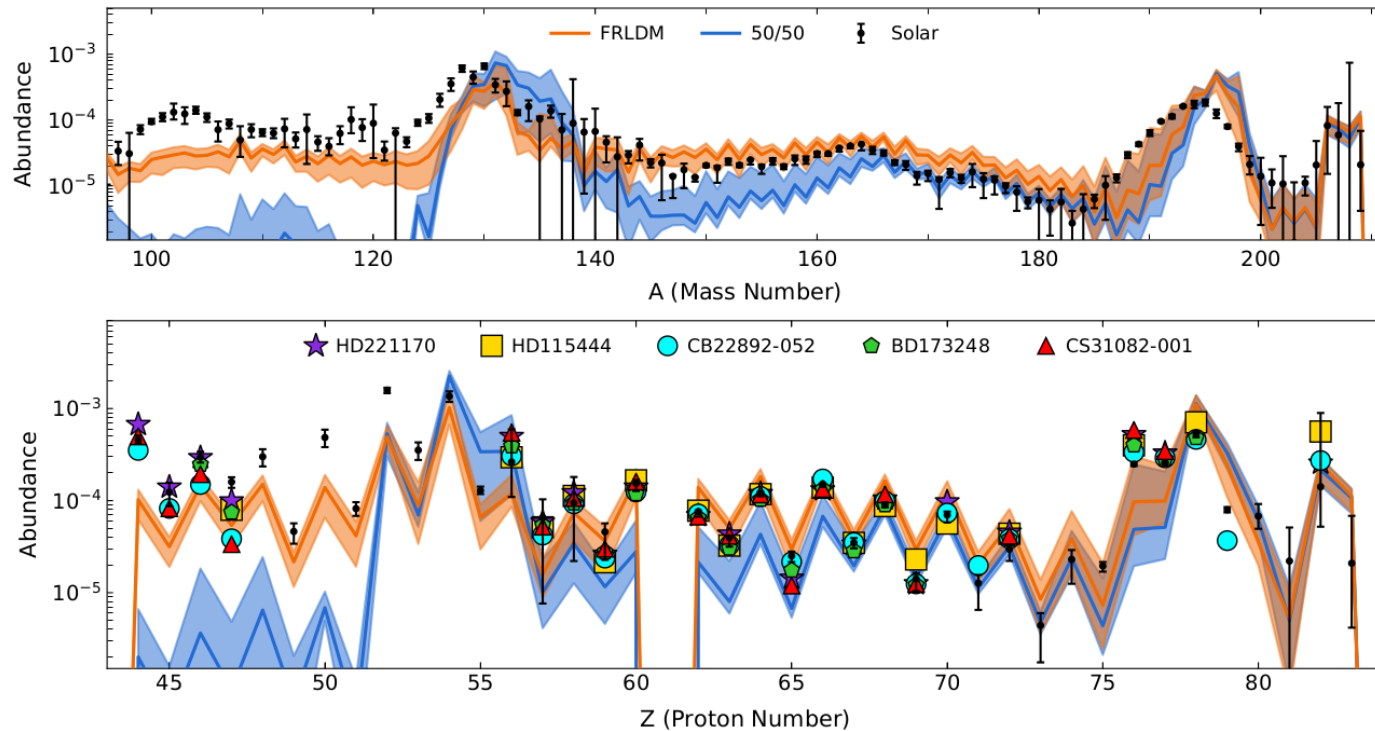


Measure the placement of material of the highest peak in  $A$

Notice the transition between 3 and 2 peaks plays an important role

$r$ -process conditions from astro. simulations suggest population of  $A \sim 150$  to higher nuclei

# IMPACT ON THE ABUNDANCES



Abundance output using **common old nuclear fission data** and **our new model predictions**

Co-production of light nuclei from  $Z \sim 45$  to the actinides (dynamical merger ejecta only!)

**Universality** may extend further down to lighter nuclei than commonly accepted in the literature

# SPECIAL THANKS TO

My collaborators

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Dolance, W. Even, C. Fontes, C. Fryer, E. Holmbeck, A.  
Hungerford, P. Jaffke, T. Kawano, O. Korobkin, J.  
Lippuner, G. C. McLaughlin, J. Miller, W. Misch, P. Möller,  
R. Orford, J. Randrup, G. Savard, T. Sprouse, R. Surman,  
N. Vassh, M. Verriere, R. Vogt, R. Wollaeger, Y. Zhu  
*& many more...*

■ Student ■ Postdoc ■ CTA Staff ■ FIRE PI



# SUMMARY

The  $r$ -process requires a *deep understanding* of fission

**Recent calculations give insight into:**

Re-cycling material ▲ Production of heaviest elements ▲ Late-time observations

FRIB, etc. will help to constrain nuclear models, but the heaviest elements will remain relatively **inaccessible**

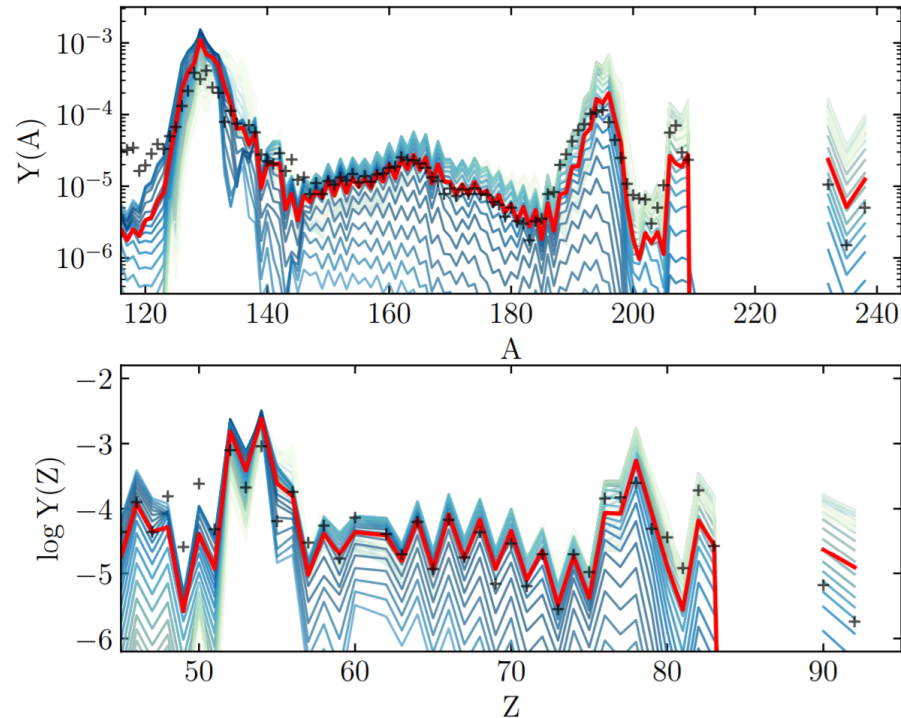
We therefore need to keep developing and studying **theoretical** models of nuclear physics, especially fission

Nuclear modeling is absolutely crucial if we want to prove definitively that heavy elements such as the actinides were made in an event

Results / Data / Papers @ [MatthewMumpower.com](https://MatthewMumpower.com)

**WHAT IS LEFT AFTER FISSION?**

# LONG-LIVED ACTINIDES

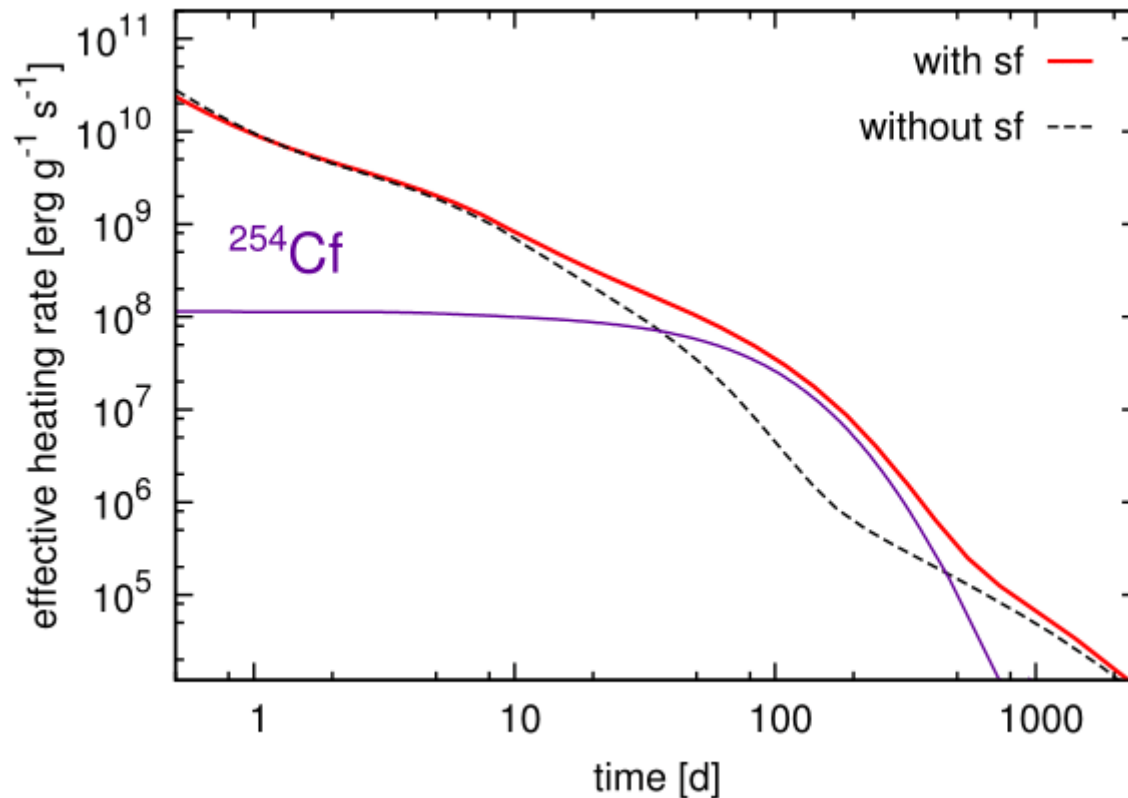


In some simulations **actinides** seem to be overproduced versus **lanthanides**

A sufficient amount of **dilution** with lighter  $r$ -process material is required to match the solar isotopic residuals

$\therefore$  Fission theory has implications far beyond nucleosynthetic outcomes; e.g. for galactic chemical evolution, etc.

# ONE EXAMPLE: $^{254}\text{Cf}$ (Z=98)

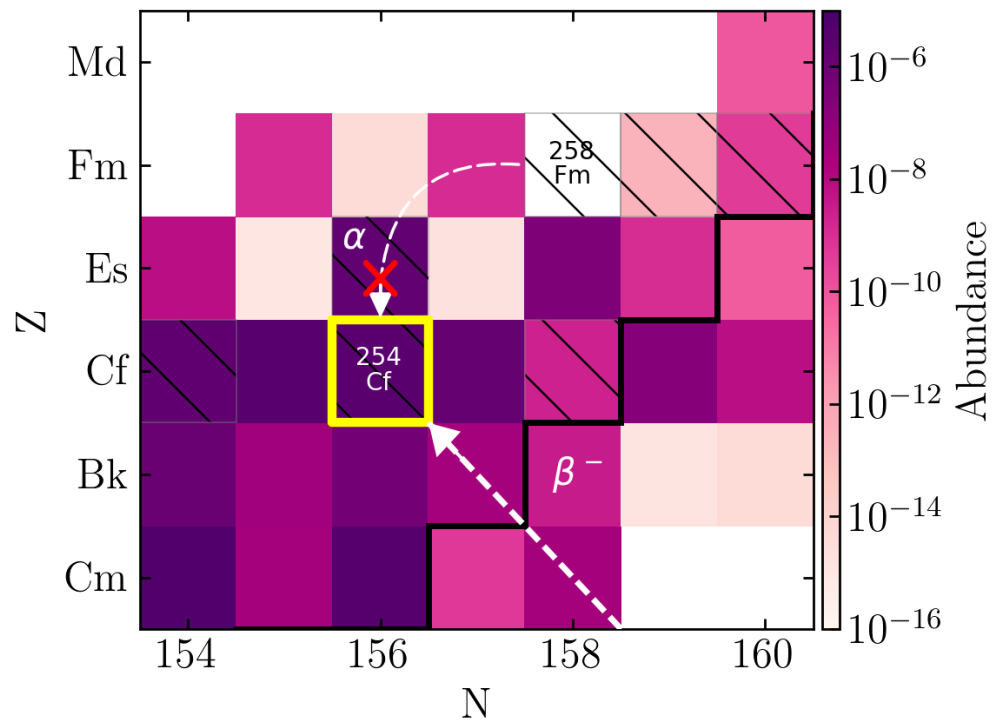
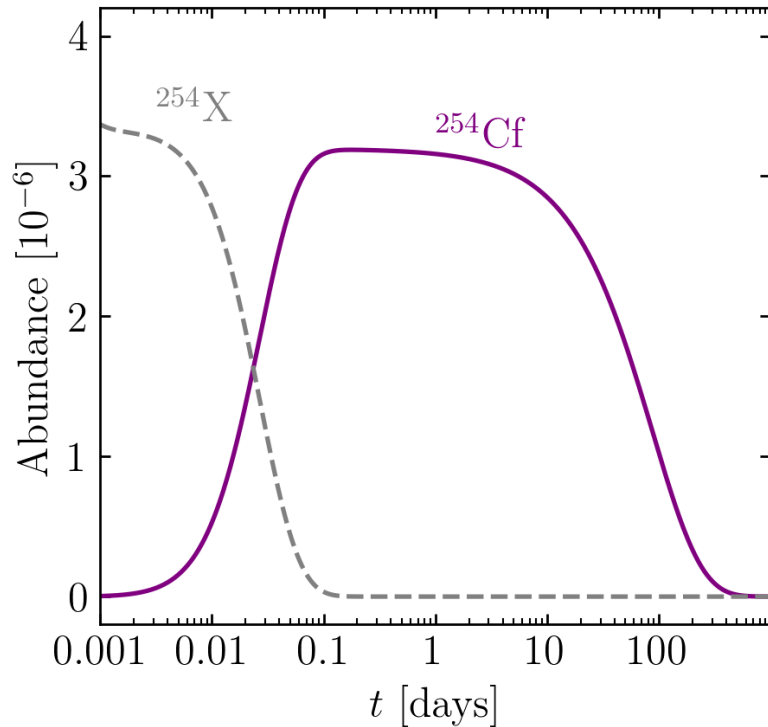


Is there any possible precursor to show that actinide nucleosynthesis has occurred in an event?... Maybe!

The spontaneous fission of  $^{254}\text{Cf}$  can be a primary contributor to nuclear heating at late-time epochs

The  $T_{1/2} \sim 60$  days; found from nuclear weapons testing

# PRODUCTION OF $^{254}\text{Cf}$ (Z=98)

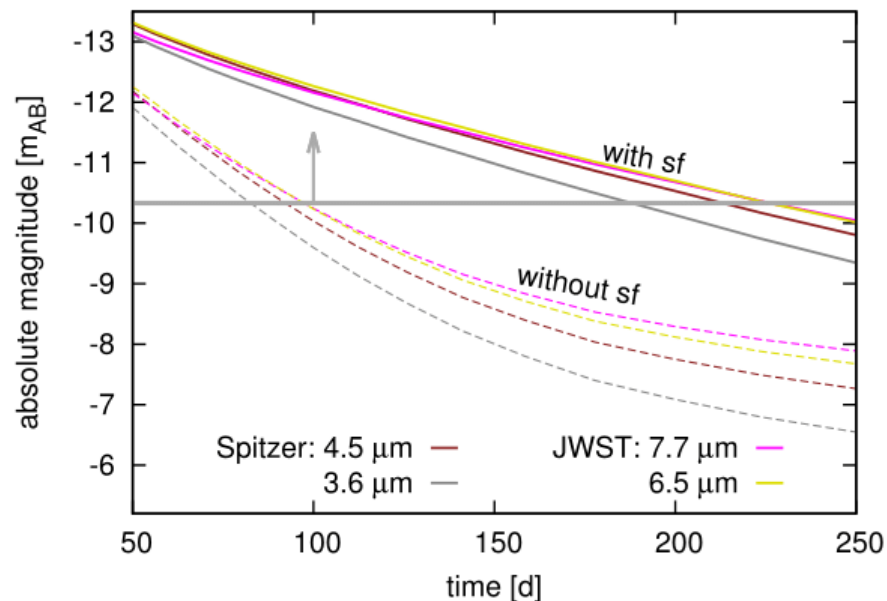
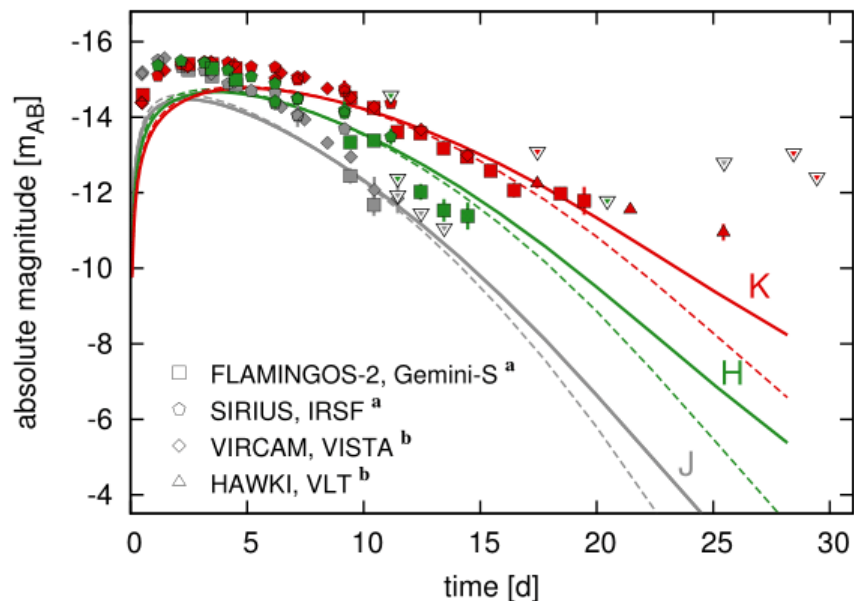


Primary feeder seems to be from  $\beta$ -decay

Production of this nucleus been explored over a range of nuclear models; some high - some low

Remains to be seen if we can disentangle from other late-time heating sources (e.g. pulsar or accretion fallback)

# OBSERVATIONAL IMPACT OF CALIFORNIUM



Both near- and middle- IR are impacted by the presence of  $^{254}\text{Cf}$

Late-time epoch **brightness** can be used as a **proxy** for **actinide** nucleosynthesis

Future JWST will be detectable out to 250 days with the presence of  $^{254}\text{Cf}$

This also has implications for merger morphology...